

Chair of Polymer Processing

## Master's Thesis

Investigating Hydrogen Diffusion in Filled Polymers with Simple Geometries: An ABAQUS Simulation Study

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### Abstract

This work was carried out in the context of the module "Polymers 4 Hydrogen" at the Polymer Competence Center Leoben GmbH. The aim of this thesis was to create a Finite Element Method (FEM) model for the simulation of hydrogen diffusion through particle filled membranes with an interface zone around the filler particles. This FEM model was based on an extended Nielsen model and was implemented in ABAQUS. At the end of this thesis, a comparison was made between the interface model, a standard FEM model without interface zone and the analytical solution of the extended Nielsen model.

For all simulations performed in this thesis, it was assumed that the polymer matrix is a homogeneous material and the filler particles act as absolute barriers which were therefore implemented as holes in the matrix. Each filler particle was modelled with a thin interface around its boundary edges that separates the particle from the matrix. In this interface zone, an orientation is applied to the mesh nodes that allows the diffusivity of the interface zone to be changed depending on the direction of flow and the adhesion coefficient. The purpose of this interface zone was to represent the true interfacial diffusion behavior between a filler particle and the matrix material.

Several ways to implement such an interface zone in an FEM model were evaluated. In the final version of the model, the interface zone was implemented at the mesh node level through the use of ABAQUS subroutines. In these models, the filler particles were regularly and periodically arranged in the membrane according to the assumptions of the Nielsen model. The evaluation and comparison of the analytical model results with the results of the FEM simulations with and without the interface zone showed that the FEM simulations with an interface zone were in better agreement with the analytical data than the simulations without the zone.

## Kurzfassung

Diese Arbeit wurde im Rahmen des Moduls "Polymers 4 Hydrogen" am Polymer Competence Center Leoben GmbH durchgeführt. Ziel dieser Arbeit war die Erstellung eines FEM-Modells zur Simulation der Wasserstoffdiffusion durch partikelgefüllte Membranen mit einer Interface-Zone um die Partikel. Dieses FEM-Modell basiert auf einem erweiterten Nielsen-Modell und wurde in ABAQUS implementiert. Am Ende dieser Arbeit wurde ein Vergleich zwischen dem Interface-Modell, einem Standard-FEM-Modell ohne Interface-Zone und der analytischen Lösung des erweiterten Nielsen-Modells durchgeführt.

Für alle in dieser Arbeit durchgeführten Simulationen wurde angenommen, dass die Polymermatrix ein homogenes Material ist und die Füllstoffpartikel als absolute Barrieren wirken, welche daher als Löcher in der Matrix implementiert wurden. Jeder Füllstoffpartikel wurde an seinen Rändern mit einer dünnen Interface-Zone modelliert, die das Partikel von der Matrix trennt. In dieser Interface-Zone wird eine Materialorientierung auf die Netzknoten angewendet, die es ermöglicht, die Diffusivität der Interface-Zone in Abhängigkeit von der Flussrichtung und dem Adhäsionskoeffizienten zu bestimmen. Der Zweck dieser Interface-Zone ist es, das tatsächliche Grenzflächendiffusionsverhalten zwischen einem Füllstoffpartikel und dem Matrixmaterial darzustellen.

Es wurden mehrere Möglichkeiten zur Implementierung einer solchen Interface-Zone in ein FEM-Modell untersucht. In der endgültigen Version des Modells wurde die Interface-Zone auf der Netzknotenebene mit Hilfe von ABAQUS-Subroutinen implementiert. In diesen Modellen wurden die Füllstoffpartikel nach den Annahmen des Nielsen-Modells regelmäßig und periodisch in der Membran angeordnet. Die Auswertung und der Vergleich der Ergebnisse des analytischen Modells mit den Ergebnissen der FEM-Simulationen mit und ohne Interface-Zone zeigten, dass die FEM-Simulationen mit Interface-Zone besser mit den analytischen Daten übereinstimmen als die Simulationen ohne Interface-Zone.

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## **Chapter 1**

## Introduction

## 1.1 Motivation

Applications in mobility and transport such as hydrogen powered vehicles make light-weight tank solutions necessary, which consist partially or completely of fiber-reinforced polymers [39]. Hydrogen, as a small molecule, easily permeates comparably dense materials. This is even more the case if the molecular structure of the dense material is not latticed, as is the case in polymers. Consequently, a polymer hydrogen tank will lose its content faster than a hydrogen tank made from metal. Therefore a reinforcement is sought to increase the diffusion barrier properties of polymers. [20]

While comparatively simple permeation models are sufficient to describe permeation in homogeneous materials, complex material systems with e.g. voids, different crystal structures, fillers, fibers, different layers, etc. are not fully understood yet and can for this reason only be modeled with simplifications. Therefore, a simulation model which is specifically designed for a more complex material system could help with predicting the diffusion processes around fillers and fibers as barrier materials in polymeric membranes.

The objective of this work is to develop a new numerical model for the simulation of diffusion through fiber-reinforced polymer composites using the Finite Element Method (FEM) solver ABAQUS [42]. The model will be built from three different phases:

• The matrix, the homogeneous polymer volume of the membrane in which diffusion occurs according to the Fickian laws and independently of direction.

- The filler particles, which are embedded in the matrix and are assumed to have ideal barrier properties. Consequently, no mass flow can occur into or out of the filler.
- The interface between the filler and the matrix. The dimension of this interface is one less than that of the entire filler-matrix geometry. This filler-matrix-interface (FMI) does not alter mass flow along its thickness, but the diffusion rate parallel to the outer edge of the filler will be assumed to be dependent on the adhesion coefficient between filler and matrix.

With this model, the barrier properties of the simulated composite can be correlated to the mechanical adhesion between the fillers and the matrix. Therefore, this model should provide a better understanding of diffusion processes in reinforced polymers.

### **1.2** State of the Art

The fundamentals of permeation and diffusion in polymers have been known for decades as shown by Crank [3], Klopffer et al. [18] and Philibert [36]. The resulting analytical models derived are still in use today and have benefited greatly from modern computing power as shown by Macher et al. [25]. However, with the increasing use of reinforced polymer composites in all areas of engineering, the need for a suitable simulation model for diffusion and permeation has arisen especially in the fields of transportation and energy. For highly simplified composite models, in which each phase is assumed to be homogeneous and all phases are separated from each other, the established models are still usable, as shown by Macher et al. [24] and Monsalve-Bravo et al. [29]. However, for more advanced and detailed simulations of filler or fiber-reinforced composites, these models must be revised. This is shown by Schultheiss [41] in his work where he used the analytical models based on Fick's laws and extended them so that they could better describe the properties of an interface layer. He uses a custom made finite difference method (FDM) solver to simulate the barrier properties of a three-phase fiber-reinforced polymer model.

Duncan et al. [5] give an overview about the basics of diffusion and permeation processes of hydrogen in polymers and describe how to model these processes. They also roughly deal with the topic of FEM simulation of diffusion processes. Furthermore, Alhijazi et al. [1] wrote a review about the possibilities of FEM analysis for natural fiber composites.

There are a few papers in the literature that use ABAQUS to simulate hydrogen diffusion through metals. Most of them use ABAQUS to simulate the mechanical weakening effects

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of hydrogen diffusion through metals. These works used the heat transfer / mass diffusion analogy and included user-defined subroutines to couple the diffusion process with the weakening of the material, e.g. [2, 4, 17, 31]. Olden et al. [32] simulated the diffusion of hydrogen through duplex stainless steels. Their model was built on two different phases (austenite and ferrite) and was used in three different setups: fine, coarse, and elongated phases. In addition, the effects of strain and stress near an embedded flaw on the hydrogen concentration were evaluated using a user-defined subroutine. Similarly, in the work of Zhang et al. [47], the effects of tensile and compressive stresses on the diffusion of hydrogen through steel pipes were simulated. A multi-phase ABAQUS simulation was performed by Pu et al. [37], on the topic of stress-assisted diffusion process along grain boundaries and the mechanical response of the grain boundary in a general polycrystalline material. Their model consists of seven different crystalline zones and the grain boundary interfaces between them.

In most of the reviewed literature references, the FEM models for the simulation of diffusion are built from only one phase and the composite behavior is added by using experimentally determined parameters, e.g. [7, 8, 13, 19, 34]. In contrast, Gholami et al. [11] used ABAQUS to simulate the hygrothermal degradation of the mechanical properties of fiber reinforced composites by performing a micro-scale analysis on a two-phase model (matrix and fibers). Li et al. [21] and Zhao et al. [48], simulated the diffusion of chloride ions through cement and they used three-phase and five-phase models, respectively. Papathanasiou and Tsiantis [35] dealt with the barrier properties of flake-filled polymers with a two-phase model in OpenFOAM [33].

### **1.3** Contribution of this Work

The novelty of this work is based on the following points:

- 1. Few papers have been found in the literature that focus on the use of ABAQUS or other commercial FEM solvers to simulate diffusion processes in reinforced polymers.
- 2. In almost all of the literature found, simplifications were made in the models used in one way or another, e.g. [7, 8, 11, 13, 19, 34, 35]. Most of the simulation models consist of only one material whose permeation parameters were determined in advance in experiments. Some simulations use a material model with two phases but these phases are depicted as homogeneous and separated from each other. These models are described in more detail in Section 1.2. The simulation model proposed in this work uses a three-phase model, which connects the matrix and the filler with the FMI.

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3. The proposed adhesion coefficient could allow more accurate modeling of complex reinforced polymer components. A model with this coefficient was not found in any of the literature reviewed. With this coefficient, it would be possible to translate the real filler matrix adhesion into an FEM model for permeation simulation. This could lead to a better understanding of the barrier properties of reinforced polymers.

## Chapter 2

## Theory

## 2.1 Permeation

Permeation is a process in which the equalization of concentration differences takes place without external influences, resulting in an increase in entropy. This physical process arises from the undirected random movements of particles as result of their thermal energy. If there are different particle concentrations in an area, these random particle movements tend to cause particles to move from higher concentrations to lower concentrations with higher probability. This results in a macroscopic mass transport which leads to an equalization of the particle concentrations.

Since all the simulations performed in this thesis are purely two-dimensional, the following equations and coefficients are also assumed to be two-dimensional. The random movements of the particles are mathematically represented by the diffusion coefficient D (m<sup>2</sup> · s<sup>-1</sup>). This coefficient comes from the Einstein-Smoluchowski relationship [6, 45] which links D with the mobility of the particles  $\mu$  (s · kg<sup>-1</sup>):

$$D = \mu \cdot k_b \cdot T \tag{2.1}$$

where  $k_b$  (J · K<sup>-1</sup>) is the Boltzmann constant and *T* (K) the absolute temperature. Another important factor is the solubility coefficient *S* (mol · m<sup>-2</sup> · Pa<sup>-1</sup>) which establishes the relationship between the concentration *C* (mol · m<sup>-2</sup>) and partial pressure *p* (Pa) of a gas in the polymer membrane:

$$C = S(T, p) \cdot p \tag{2.2}$$

*S* can be dependent on *T*, *p* or both. If the diffusion coefficient *D* is multiplied with the solubility coefficient *S* the permeation coefficient *P* (mol  $\cdot$  s<sup>-1</sup>  $\cdot$  Pa<sup>-1</sup>) is obtained:

$$P = D \cdot S \tag{2.3}$$

This coefficient is the link between the kinetic and thermodynamic aspects of diffusion and is a measure of how well a gas can permeate through a solid at a given pressure. [10, 18]

## 2.2 Sorption Models

In Table 2.1, five established sorption models are listed. The first three are explained in detail below as they are relevant for the diffusion of gases through polymer membranes. The last two models are mainly applied on the diffusion of vapors which tend to condense, and therefore cause swelling in polymers, and are only listed for the sake of completeness.

Table 2.1 Different models of sorption and typical associated interactions. [18]

Sorption model	Main component interactions	
Henry	polymer-polymer	
Langmuir	polymer-penetrant	
Dual mode	combination of Henry and Langmuir models	
Flory-Huggins	penetrant-penetrant	
BET	combination of Langmuir and Flory-Huggins models	

• Henry's law sorption: This is the simplest sorption case where the gas can be assumed to be ideal and the relationship between the concentration and the pressure of the gas can be described as linear. This model works best at low pressures where the polymer-polymer interactions are stronger than the polymer-penetrant or penetrant-penetrant interactions. The solubility coefficient *S* (Eq. 2.2) is constant for this model. Thereby, the equation for Henry's law can be written as followed: [18]

$$C = S \cdot p \tag{2.4}$$

• Langmuir mode sorption: In this sorption model, a linear relationship between pressure and concentration in the membrane at low pressures is assumed. Above a certain

concentration, this sorption model depicts a saturation of the membrane. Once the saturation is reached, the concentration in the membrane will stay constant, even if the pressure is increased further. The saturation concentration of the membrane is described by  $C'_H$  (mol·m<sup>-2</sup>) and  $b_H$  (Pa<sup>-1</sup>) describes the gradient of the linear segment at the start of the model.

$$C = \frac{C'_H \cdot b_H \cdot p}{1 + b_H \cdot p} = S(p) \cdot p \tag{2.5}$$

• Dual mode sorption: The dual mode model is a combination of Henry's law model and the Langmuir model. It was developed to describe the sorption of non-reactive gases in glassy polymers. This model considers the presence of two different diffusing molecule types: one trapped and one freely moving. [18]

In this work, the Henry's law sorption model is applied, because it is assumed that the simulations operate in a region in which a linear relationship between gas pressure and concentration is valid. Furthermore, hydrogen can be assumed to behave ideally at such low pressures and the membrane has no trapping sites.

### **2.3 Basic Diffusion Models**

In 1855, Adolf Fick empirically derived the two basic laws of diffusion, the so-called Fick's laws. These were theoretically derived from thermodynamics, and thereby proven, by Albert Einstein at the beginning of the 20th century. [6]

Fick's first law describes the relationship between the diffusive flow density  $F \pmod{m^{-1} \cdot s^{-1}}$ and the gradient of the concentration  $\nabla C \pmod{m^{-3}}$ . In this work, it is assumed that *D* is independent of *C* and is therefore not derived. This assumption is reasoned in Section 2.4. This results in the equation for Fick's first law:

$$F = -D \cdot \nabla C \tag{2.6}$$

Fick's second law is described by a partial differential equation relating the change in concentration over time  $\frac{\partial C}{\partial t}$  (mol · m<sup>-2</sup> · s<sup>-1</sup>) to the Laplace operator of the concentration *C*. This results in the following formula for Fick's second law:

$$\frac{\partial C}{\partial t} = D \cdot \Delta C \tag{2.7}$$

## 2.3.1 Analytical Solution of Fick's Laws for Diffusion through a Membrane

This section gives a brief overview of the analytical simulation of diffusion. For a detailed description of the mathematical derivation of the equation shown below, see Crank [3] or Macher et al. [25].

At the start of the permeation process the initial state is in effect. There is no concentration of the permeate in the membrane and the diffusion process is just beginning. As soon as the diffusion process has started, it is in the transient state. Here the concentration of the permeate in the membrane increases with time. The permeate is deposited in the membrane, so to speak. The saturation of the increase in concentration of the permeate in the membrane depends on the boundary conditions of the membrane. For example, if the membrane is immersed in the medium. Then each side of the membrane would be exposed to the same level of medium concentration  $C_{sub}$ . In this case, saturation of the membrane is reached when  $C_{sub}$  is present at every point in the membrane. Another case would be that the membrane acts as a barrier of some kind. Then on one side of the membrane the concentration of the permeate would be high  $(C_{high})$  and on the other side of the membrane the concentration of the permeate would be low  $(C_{low})$ . In this case, saturation of the membrane would be reached when a linear gradient of concentration across the membrane is reached. When the membrane has reached such a saturated state, the transient state is over and the steady state is reached. At steady state, the rate of diffusion into the membrane is equal to the rate of diffusion out of the membrane. Therefore, the concentration of the permeate in the membrane does not change in this state. This means that reaching steady state does not mean that the membrane is completely saturated, but that it has reached a saturation state according to the boundary conditions. The initial and boundary conditions shown in Eq. 2.8 allow the evaluation by an analytical method. Only if the concentration at the inlet and outlet surface remains constant, an analytical calculation of the diffusion can be derived.

$$C(x,t=0) = 0 \quad \text{for} \quad 0 < x < L$$

$$C(x=0,t) = C_1 \quad \text{for} \quad 0 < t < \infty$$

$$C(x=L,t) = C_2 \quad \text{for} \quad 0 < t < \infty$$
(2.8)

If Eq. 2.7 is solved as a Sturm-Liouville problem with the initial and boundary conditions shown in Eq. 2.8 by using eigenfunctions, the result for the concentration is

$$C(x,t) = (C_2 - C_1)\frac{x}{L} + C_1 + \frac{2}{\pi} \cdot \sum_{n=1}^{\infty} [(-1)^n \cdot C_2 - C_1] \cdot \sin(n\pi \frac{x}{L}) \cdot e^{-\frac{n^2 \pi^2 D t}{L^2}}$$
(2.9)

In order to calculate the flow out of the membrane as a function of time, the analytical solution of the concentration (Eq. 2.9) is substituted into Fick's first law (Eq. 2.6) with differentiation at x = L, so that the result is

$$F(t) = D \cdot \frac{C_1 - C_2}{L} + \frac{2D}{L} \cdot \sum_{n=1}^{\infty} \left[ (-1)^n \cdot C_1 - C_2 \right] \cdot e^{-\frac{n^2 \pi^2 D_t}{L^2}}$$
(2.10)

If Eq. 2.10 is integrated by t with limits 0 (s) and t (s), the following equation for the cumulative flow density  $Q \pmod{(\text{mol} \cdot \text{m}^{-1})}$  out of the membrane for time t is obtained:

$$Q(t) = D \cdot \frac{C_1 - C_2}{L} \cdot t - \frac{(C_1 + 2C_2)L}{6} - \frac{2L}{\pi^2} \cdot \sum_{n=1}^{\infty} \frac{[(-1)^n \cdot C_1 - C_2]}{n^2} \cdot e^{-\frac{n^2 \pi^2 Dt}{L^2}}$$
(2.11)

In general, it is advantageous for the application of an analytical solution in a numerical calculation that its quantities are changed into a dimensionless form because this reduces rounding errors and simplifies fitting algorithms, as the dimensionless equations are easily scalable. The following equations show an example of such a transformation for the concentration C, the distance x and the time t.

$$\hat{C} = \frac{C}{C_1}, \quad \hat{x} = \frac{x}{L}, \quad \hat{t} = \frac{D \cdot t}{L^2}$$
 (2.12)

where  $C_1$  is the boundary condition of the concentration at the upstream face of the membrane and L is the total thickness of the membrane.  $\hat{C}$ ,  $\hat{x}$  and  $\tau$  are the respective dimensionless variables.

### 2.4 Classification of Diffusion Modes for Polymers

Diffusion is a process at the molecular level and therefore both the type of diffusing gas and the structure of the polymer membrane are crucial. As shown in Table 2.2, the glass transition temperature  $T_g$  of the polymer membrane and the critical temperature  $T_c$  of the gas are of crucial importance.

T value compared		More condensable gases
to a characteristic	Gases with $T > T_c$	or vapors $(T < T_c)$
temperature of	H <sub>2</sub> , He, O, N <sub>2</sub>	$CO_2$ , $SO_2$ ,
the system		NH <sub>3</sub> , hydrocarbons
	Fickian diffusion	Fickian diffusion
	constant D	<i>D</i> function of $C: D(C)$
$T > T_g$	Henry's mode sorption	Single mode sorption
<b>Rubbery polymers</b>	constant <i>S</i> increases slightly with <i>T</i>	S decreases with $T$
	<i>P</i> decreases slightly with pressure	<i>P</i> increases with pressure
	(hydrostatic pressure effect)	(plasticization effect)
T < T <sub>g</sub>	Dual mode sorption $S(p)$	Dual mode sorption $S(p)$
Glassy polymers	Free volume diffusion	Non-Fickian and
		anomalous diffusion

Table 2.2 General behavior observed for the transport of small molecules in polymers.[16, 26]

Depending on the ambient temperature T, the following four cases can be distinguished [18, 40]:

- *T* is higher than  $T_c$  and  $T_g$ : This case is most easily represented analytically or by computer simulation. Diffusion proceeds strictly according to Fick and it is reasonable to calculate with constants *D*, *S* and *P* values. Henry's mode sorption applies, whereby the pressure and concentration of the gas have a linear relationship at low gas pressures. This model reaches its limits if the compressibility of the gas has a significant influence.
- *T* is lower than  $T_c$  and higher than  $T_g$ : This case can still be described by purely using Fick's equations. As a rule, *S* can still be assumed to be constant. On the other hand, *D* shows a dependence on *C*, which means that *P* can also be assumed to be variable.
- *T* is higher than  $T_c$  and lower than  $T_g$ : In this case, Fick's laws can no longer be applied; instead, free volume diffusion methods are used. Furthermore, dual mode diffusion can be assumed in this case. This means that there is a Henry diffusive flux and a Langmuir diffusive flux. The resulting two diffusion coefficients can be combined

into a p-dependent diffusion coefficient. S is also p dependent, but S and D are not interdependent. P can also be assumed to be variable.

• T is lower than  $T_c$  and  $T_g$ : As in the previous case, all three coefficients are variable. However, this area is absolutely Non-Fickian and the methods of abnormal diffusion are used.

As mentioned earlier in Section 2.2, Henrys' law sorption and Fickian diffusion are assumed for this work. Furthermore, this work focuses on the diffusion of hydrogen gas. Therefore, of the four cases shown in Table 2.2, only the upper left is applicable to this work.

### 2.5 Diffusion in Particle Filled Polymer Systems

The focus of this work is on particle-filled polymer membranes. Therefore, this section serves as a brief introduction to the theoretical models that have been formulated for this type of multi-component material system. An overview of the structure of such a filler-polymer system is given in Fig. 2.1. The particles shown in this figure are regularly and periodically arranged as described by Nielsen in his work [30]. Furthermore, the distances between the fillers and the rows of fillers are defined by *s* and *d* as described by Macher et. al. [24] in their work, which serves as an extension of the model proposed by Nielsen. This extension of the model also allows it to be implemented in numerical simulations. This in turn allows a direct comparison between the analytical and numerical data obtained.

#### 2.5.1 Nielsen Model

Fick's First law (Eq. 2.6) is used to calculate the flow of the permeate. In this work, it is assumed that the sorption behaves according to Henry's law (Eq. 2.4). By combining these two laws, the following equation is obtained:

$$F = -D \cdot S \cdot \frac{dp}{dx} \tag{2.13}$$

with  $\frac{dp}{dx}$  (Pa · m<sup>-1</sup>) as the pressure gradient along the thickness of the membrane. Furthermore, since this work assumes steady-state behavior for all diffusion processes, the following equation can be derived.



**Fig. 2.1** Schematic drawing of a particle-filled polymer membrane, F represents the diffusion flow density and the arrow indicates the direction, w (m) is the width of a filler particle perpendicular to F and b (m) is the thickness of a filler particle parallel to F, s (m) is the slit shape which describes the distance between two filler particles in a row and d (m) is the filler distance which describes the distance from one row of filler particles to another. [24]

$$F = -D \cdot S \cdot \frac{p_1 - p_2}{L} \tag{2.14}$$

Where  $p_1$  (Pa) and  $p_2$  (Pa) are the pressures on the side of the membrane with the high and low concentration of the permeate, respectively, and L (m) is the thickness of the membrane between these two sides. Nielsen described filler particles as impermeable barriers in the membrane [30]. With this assumption, Eq. 2.15 describes the effective length  $L_{eff}$  (m) that the permeate has to travel through the membrane.  $\tau$  (-) is the tortuosity factor that proportionally describes the longer tortuous path of the permeate due to the impermeable filler particle assumption.

$$L_{eff} = \tau \cdot L \tag{2.15}$$

On average, a filler particle in the membrane can be expected to add  $\frac{w}{2}$  to the length of  $L_{eff}$  [30]. Each diffusing particle encounters on average a number of filler particles  $\langle N \rangle$ , which is calculated as described in Eq. 2.16. *b* is the thickness of the filler particles as shown in Fig. 2.1 and  $\phi_f$  is the volume fraction of the filler particles in the membrane.

$$\langle N \rangle = \frac{L}{b} \cdot \phi_f \tag{2.16}$$

With these two assumptions, the effective length of the path of the permeate can be calculated as

$$L_{eff} = L \cdot \left(1 + \frac{w}{2b} \cdot \phi_f\right) \tag{2.17}$$

Where w is the width of the filler particle as shown in Fig. 2.1. By comparing Eq. 2.15 and Eq. 2.17, the tortuosity factor of the filled membrane can be described as

$$\tau = 1 + \frac{\alpha}{2} \cdot \phi_f \tag{2.18}$$

with  $\alpha$  being the aspect ratio of the filler particles defined as  $\frac{w}{h}$ .

For convenience, it is suggested not to change the external measurable parameters such as L,  $p_1$  or  $p_2$ . Instead, an effective diffusion coefficient  $D_{eff}$  can be defined according to Eq. 2.14 and Eq. 2.17:

$$D_{eff} = \frac{D_0}{1 + \frac{\alpha}{2} \cdot \phi_f} \tag{2.19}$$

where  $D_0$  is the diffusion coefficient of an unfilled membrane.

Furthermore, since the filler particles are assumed to be impermeable, they reduce the available volume for permeation. The effective solubility coefficient is calculated in Eq. 2.20, in which  $S_0$  is the solubility coefficient of an unfilled membrane.

$$S_{eff} = S_0 \cdot (1 - \phi_f) \tag{2.20}$$

In essence, the Nielsen model can be written as the ratio of the effective flow  $F_{eff}$  through a filled membrane to the flow  $F_0$  through an unfilled membrane of the same external dimensions. By combining this statement with Eqs. 2.14, 2.19 and 2.20, the following equation can be derived.

$$\frac{F_{eff}}{F_0} = \frac{D_{eff} \cdot S_{eff}}{D_0 \cdot S_0} = \frac{P_{eff}}{P_0} = \frac{1 - \phi_f}{1 + \frac{\alpha}{2} \cdot \phi_f}$$
(2.21)

#### 2.5.2 Extension of the Nielsen Model

As shown in Eq. 2.21, the Nielsen model depends on only two variables,  $\alpha$  and  $\phi_f$ . As mentioned above, this is not enough information about the particle filled membrane to model a numerical simulation with it. An additional specification of the geometric dimensions of the filled membrane is necessary to directly compare Nielsen's model with the results of a numerical simulation. As shown in Fig. 2.1, Macher's extension of Nielsen's model introduces the slit shape *s* and the filler distance *d*. With these two additional variables, it is possible to reproduce the same model in a numerical simulation. Furthermore, this figure shows that there are channels with unobstructed flow paths between the filler particles, which was also predicted by Nielsen in his work [30]. This means that the flow of the permeate through the filled membrane consists of two partial flows. One is the unobstructed flow through the channel and the other is the tortuous flow that is obstructed by parts of the fillers. This configuration of the filled membrane can be reduced to a unit cell model shown in Fig. 2.2. There is also a representation of the paths of the two types of flow through the filled membrane. This approach is similar to the one used by Minelli et. al. [27] in their work.



**Fig. 2.2** Schematic representation of a unit cell in a filled membrane structured as described by Nielsen [30]. The unit cell is enclosed by the dashed lines. (a) Dimensions of the components in the unit cell. The blue area represents a channel with a width of  $w_{ch}$  (m) where the permeate is not hindered in its flow. (b) The two accumulating flows through the unit cell are shown.  $F_{ch}$  in blue is the unobstructed flow as described in (a) and  $F_t$  in red is the tortuous flow around the filler particles. [24]

The area  $A_u$  (m<sup>2</sup>) of such a unit cell can then be calculated as shown in Eq. 2.22. Furthermore, the area  $A_f$  (m<sup>2</sup>) of the filler particle in this unit cell is then defined according to Eq. 2.23.

$$A_u = \frac{w+s}{2} \cdot d \tag{2.22}$$

$$A_f = \frac{w \cdot b}{2} \tag{2.23}$$

Due to the two-dimensionality of this model,  $\phi_f$  (-) is reintroduced as the filler area fraction, which can be calculated as

$$\phi_f = \frac{A_f}{A_u} = \frac{w \cdot b}{d \cdot (w+s)} \tag{2.24}$$

By rearranging Eq. 2.24, the filler distance can then be calculated as

$$d = \frac{w \cdot b}{\phi_f \cdot (w+s)} \tag{2.25}$$

As shown in Eq. 2.25, the filler distance depends on the slit shape *s* of the filled membrane.

To calculate *s*, it is necessary to determine the ratio between the flows through a filled and an unfilled membrane. As shown in Fig. 2.2, the flow through the filled membrane consists of the two flows  $F_{ch}$  and  $F_t$ . Assuming that the channels in the filled membrane are constant along the thickness of the membrane, the flow through them depends only on their width. The width of these channels can be calculated as

$$w_{ch} = \frac{s - w}{2} \tag{2.26}$$

If the width of the channels in the filled membrane is related to the width of the whole unit cell, the channel flow ratio  $f_{ch}$  (-) can be calculated as the fraction of the flow through the channel  $F_{ch}$  and the flow through the unfilled membrane  $F_0$ .

$$f_{ch} = \frac{F_{ch}}{F_0} = \frac{s - w}{w + s}$$
(2.27)

To calculate the tortuous flow  $F_t$ , it is necessary to calculate the effective length  $L_t$  (m) of this flow through the unit cell as shown in Eq. 2.28. The solution of the integral in this equation results in the average additional path length x (m) caused by the filler particle in the unit cell.

$$L_t = d + \frac{\int_0^{\frac{w}{2}} x \, dx}{\frac{w}{2}} = d + \frac{w}{4} \tag{2.28}$$

#### 2 Theory

By combining Eq. 2.15 with Eq. 2.28, the following equation can be derived for the tortuosity of the effective path length.

$$\tau_t = 1 + \frac{w}{4d} \tag{2.29}$$

Before the tortuous ratio  $f_t$  (-) of the flow through the filled membrane can be calculated, another assumption must be made. As shown in Fig. 2.2, the tortuous flow is not calculated per unit cell, but per two unit cells connected in the direction of flow. With this arrangement of two unit cells it is possible to calculate the tortuous flow ratio with only one equation:

$$f_t = \frac{F_t}{F_0} = \frac{1}{2\tau} \cdot \frac{w}{w+s} \tag{2.30}$$

Combining Eq. 2.30 with Eqs. 2.25 and 2.29 results in the final equation for the tortuous ratio of flow through the filled membrane as

$$f_t = \frac{2w \cdot b}{\phi_f \cdot (w+s)^2 + 4b \cdot (w+s)}$$
(2.31)

The sum of  $f_{ch}$  and  $f_t$  is equal to the relative flow through the filled membrane  $\frac{F_{eff}}{F_0}$  and can be substituted into Eq. 2.21. This results in

$$\frac{F_{eff}}{F_0} = \frac{1 - \phi_f}{1 + \frac{\alpha}{2} \cdot \phi_f} = f_{ch} + f_t \tag{2.32}$$

Substituting Eqs. 2.27 and 2.31 in Eq. 2.32 defines the following equation for the slit shape.

$$s = w \cdot \left[ -\frac{2\alpha \cdot \phi_f + 4}{\phi_f \cdot \alpha \cdot (\alpha + 2)} + \frac{\sqrt{\alpha^4 \cdot \phi_f^2 + 6\alpha^3 \cdot \phi_f + 12\alpha^2 + 24\alpha + 16}}{\phi_f \cdot \alpha \cdot (\alpha + 2)} \right]$$
(2.33)

With Eqs. 2.25 and 2.33 in place, the missing parameters s and d of the Nielsen model are determined and the now extended model can be compared to numerical simulations. It is important that one of the following two statements in Eq. 2.34 is true so that there is no collision between filler particles.

$$s > w$$
  
or (2.34)  
 $d > b$ 

## 2.6 Interface Layer

Every multi-component material contains interfacial layers formed by chemical and physical processes. These processes include interdiffusion of atoms or molecules, immobilization, crystallization of thermoplastics, and crosslinking. The interfacial layer is a thin transition zone between the matrix and a filler, with a thickness in the nanometer range. The structural, physical, chemical and mechanical properties of the interfacial layer may differ from those of the matrix and filler, depending on the type of filler and matrix and the physical and chemical processes involved. The existence of this interfacial layer has been proven by Fourier transform infrared (FTIR) spectroscopy [12]. Each interfacial zone has different properties due to the many influences acting on it. In general, the interface is considered the weakest part of a multi-component material, but it also provides the opportunity to tailor the material to provide the required properties [16].

Because of their variability, but also because of the numerous factors that influence their properties, there have been many studies on interface layers. A lot of research has been done on carbon fibers due to their popularity in the industry, (e.g. [9, 22, 23, 46, 49]). But there is also literature for various other material combinations such as polymer-natural fiber, polymer blends, nanocomposites, etc. [14, 28, 38, 43, 50].

## **Chapter 3**

## Simulation

For this work, several simulations were carried out for three different simulation models. The specifyers of these models are:

- Membrane with a circular filler particle, shown in Fig. 3.1a
- Membrane with a rectangular filler particle, shown in Fig. 3.1b
- Membrane with multiple rectangular filler particles, shown in Fig. 3.1c

The first two models, each with just a single filler particle in the center, were primarily used to enable the development of the interface zone simulation technique. This technique is described in more detail in Section 3.1. These simple models were also used to gather data for parameters such as interface zone thickness, interface zone mesh size, et cetera. These parameters were necessary for the successful implementation of the interface zone in the more complex third model. Several tests were carried out with this model using different filler particle aspect ratios and filler particle area ratios. These were then compared with the results of the same simulations without the interface zone and the results of the extended Nielsen model, which was described in Section 2.5.2.

Each model in this work consists of a two-dimensional membrane containing one or more two-dimensional filler particles. The filler particles are either circular or rectangular in shape and have an interface zone and a mesh transition zone around them. A simplified schematic of such a membrane with a rectangular filler particle in the center is shown in Fig. 3.2. The purpose of the interface zone is to create a thin layer around the filler particle that increases or decreases the diffusion rate of the membrane along the filler edge. Therefore, the diffusion rate normal to the filler particle edge remains unchanged from the diffusion



**Fig. 3.1** Geometries of membranes with a) a circular, b) a rectangular and c) multiple rectangular filler particles. The inlets of these membranes are always at the top edge, the outlets are always at the bottom edge, and the left and right edges of the membranes are connected by a periodic boundary condition.

rate of the membrane and only along the edges of the filler a change is applied. The mesh transition zone is only used to provide a more gradual transition of the mesh for a stable mesh generation process.

### **3.1 Interface Zone**

The core of this work is the implementation of the interface zone. Therefore, the process leading to the final configuration of the interface zone is discussed in detail in this section. As described in Section 2.6, a real membrane, reinforced with filler particles would show a change in the properties of the membrane in the transition area from filler to membrane. This effect occurs most strongly directly at the edge of the filler and bleeds into the membrane material. Since the significant portion of the effect occurs directly at the filler edge, it was initially decided to define the interface zone as a one-dimension lower element, in comparison to the rest of the model. That is, if the geometry of the membrane and the filler particle are assumed to be two-dimensional, the interface zone would have had to be implemented as a



**Fig. 3.2** Simple overview of the membrane with a rectangular filler particle, the mesh transition zone and the interface zone.

one-dimensional element. This would have allowed for the simplest possible model setup and kept the simulation time low. In ABAQUS, however, it is not possible to implement a one-dimensional element, i.e. an edge, with its own material properties in a two-dimensional diffusion analysis. There are so-called STRINGER elements in ABAQUS which allow to assign a material to an edge, but these also require the specification of a cross-sectional area for this edge. This means that the one-dimensional behavior is only superficial, and behind a black box the simulation operates in two dimensions. This was unacceptable because it would lead to results that were difficult or impossible to evaluate. Therefore, it was decided to implement a custom two-dimensional solution that could best represent the real behavior of such an interface zone. This process would eliminate the black box and allow more control over the simulation process.

The first approach in this variant of implementation was to partition the membrane in such a way that a separate area is created around the filler particle. This area was then assigned its own direction-dependent material. Furthermore, the local orientation of the coordinate system in this area was rotated in such a way that the above-described effect of a changed flow around the filler was created. If the thickness of this area was small enough, the results of these simulations were a step in the right direction. However, due to the sharp transition

of the diffusion coefficients at the boundaries of the interface zone, undesirable edge effects occurred which significantly degraded the simulation quality.

In order to keep the interface zone as thin as possible and still ensure that the materials show a gradual transition, it was decided to change the material properties at the mesh node level. For this purpose it was necessary to integrate ABAQUS subroutines into the simulation. ABAQUS provides more than 100 different subroutines for different applications. Each of these subroutines has access to different systems or data from ABAQUS and has its own methods to interact with ABAQUS. To achieve the desired effect, three different subroutines, ORIENT, UFIELD, and USDFLD, were programmed and implemented in Fortran [15] as described in Section 3.2.

In the first iteration the handling of these subroutines was mostly unknown so the subroutines were powered by a Python [44] script which calculated the necessary field data. In this Python script, the positions and sizes of the filler particles as well as the positions of the mesh nodes were read from the ABAQUS input file. Then the distance of the nodes to the filler particles was calculated and if their distance to the particle edge was below a threshold value, they were assigned a scalar value representing the relative position between the filler particle edge and the threshold value. Furthermore, in the script the interface zone was divided into the areas shown in Fig. 3.3 and for each node the corresponding area was evaluated. This enabled the calculation of the necessary rotation of the local coordinate system for each node. The values for distance and rotation calculated in this script were then written to CSV files. Of the three subroutines used, UFIELD and USDFLD were responsible for adjusting the diffusion coefficients and ORIENT for rotating the local coordinate systems of the nodes. For this purpose, the relevant CSV file was read in each case and the correct value was imported and applied. This procedure had the disadvantage that a lot of file calls had to be made which increased the simulation time significantly. For this reason, the membrane was first divided into two materials, the interface material and the membrane material. The interface material is linked to the subroutines, i.e. the subroutines are only called for mesh nodes and mesh elements for which this material has been assigned. The membrane material is a homogeneous material with constant values that is assigned everywhere except in the interface zone. This minimized the number of subroutine calls. However, simulations with many filler particles still resulted in long simulation times due to the mesh density in the interface zone. Therefore, the Python script tasks were shifted to the subroutines, reducing the number of file calls to one per type of subroutine. The exact structure and details of the functions of the final subroutines will be discussed in the next section.



**Fig. 3.3** Simple sketch of the interfacial zone and its partitions around a rectangular filler particle. Partitions 1 and 5 change the diffusive flow along the vertical axis, partitions 3 and 7 along the horizontal axis, and partitions 2, 4, 6, and 8 always at a 45 degree angle between the two neighboring partitions.

### **3.2** Fortran Subroutines

The subroutine ORIENT is responsible for rotating the local coordinate system of each mesh node to achive the desired directional behavior of the interface zone. This means that during the simulation this subroutine is called for each mesh node in the interface zone. Two different sequences exist for this subroutine, depending on whether the filler particle is circular or rectangular. For the variant used for rectangular fillers, an example is given in Algorithm 1 in the form of pseudocode. The code for a filler particle with a circular shape differs in that it is not necessary to divide the interface zone into separate orientation partitions since the direction vector from the filler center to the node can be calculated and the local coordinate system can be rotated based on it. The final result of this subroutine is shown in Fig. 3.4. By comparing the circular and rectangular filler particles there, it is easy to see that in the case of the circular filler particle, there is a continuous change of orientation along the edge of the filler, while in the case of the rectangular filler particle, different zones are clearly visible. The actual code used for the subroutine is provided in the appendix in Section A.1.1.



**Fig. 3.4** Oriention of the loacal node coordinate system in the interface zone for a circular filler particle: (a) complete membrane and (b) zoomed view. Oriention of the loacal node coordinate system in the interface zone for a rectangular filler particle: (c) complete membrane and (d) zoomed view. In this work, the 2-axis shown in yellow is the one with the variable diffusion rate.

The subroutine UFIELD is responsible for generating a field over the interface zone around the filler particles. This field contains a scalar value for each node in the interface, which indicates the relative position of the node with respect to the thickness of the interface zone. This means that if a node is located directly on the edge of a filler particle, the value of this field for this node is equal to one. On the other hand, if a node is located directly on the outer edge of the interface zone, the field has a value of zero for this node. An example of this subroutine is given in Algorithm 2 in the form of pseudocode. The final result of this subroutine can be seen in Fig. 3.5. The actual code used for the subroutine is provided in the appendix in Section A.1.2.

Algorithm 1: Pseudocode example for an ABAQUS ORIENT subroutine. **SUBROUTINE** ORIENT /\* The data of the node for wich this subroutine is called. \*/ Input: nodeData; /\* Create a global list, which can be called from every subroutine call, in which the data of all the filler particles is stored. \*/ **Create Global List:** fillerPartikelList = EMPTY; /\* Create variables to hold numerical data. \*/ **Create Variable:** minDistanceToFillerParticle = INFINITY; /\* Create variables to hold numerical data. \*/ **Create Variable:** idClosestFillerParticle = NULL; /\* If the filler particle list is empty, load the necessary data from a csv file. This is only necessary for the first time this subroutine is called. \*/ **if** *fillerPartikelList* == *Empty* **then load** *fillerParticleData* **into** *fillerPartikelList;* /\* Search for the closest filler particle to the node and save its id. \*/ for fillerParticle in fillerPartikelList do calculate distance from node.position to fillerParticle; if distance < minDistanceToFillerParticle then minDistanceToFillerParticle = distance; idClosestFillerParticle = fillerParticle.id; /\* Evaluate the orientation zone id for the current node based on its position. \*/ calculate orientationZone; /\* Apply a rotation to the local coordinatesystem of the node based on the orientation zone. \*/ **if** *orientationZone* = = 1 **then** rotate localCoordinateSystem LEFT 90 DEGREES; else if *orientationZone* == 2 or 6 then rotate localCoordinateSystem LEFT 45 DEGREES; else if *orientationZone* == 3 or 7 then Pass: else if *orientationZone* == 4 or 8 then rotate localCoordinateSystem RIGHT 45 DEGREES; **else if** *orientationZone* == 5 **then** rotate localCoordinateSystem RIGHT 90 DEGREES;



**Fig. 3.5** Diffusion coefficient field in the interface zone for (a) a circular filler particle and (b) a rectangular filler particle.

The subroutine USDFLD is significantly different from the two previously described. This subroutine is called for all mesh elements in the interface zone and not for the mesh nodes. This is because this subroutine is integrated in the material behavior and thus directly accesses the integration points of the elements. It follows that the task of this subroutine in this work is to interpolate the values of the field generated by the UFIELD subroutine to the integration points of the mesh elements. Since this is the basic function of this subroutine, it was implemented without any additional code. It was not possible to output the result of this subroutine graphically by ABAQUS, because it is only a subfunction and not an independent field. The actual code used for the subroutine is provided in the appendix in Section A.1.3.

The field generated by the USDFLD subroutine is then used by the ABAQUS material itself to assign a diffusion coefficient to the integration points based on this exponential function:

$$D_{IP}(x) = D_{Imax} \cdot e^{\Psi \cdot (x-1)} + D_M \cdot (1 - e^{\Psi \cdot (x-1)})$$
(3.1)

In this equation,  $D_{IP}$  is the diffusivity coefficient in the interfacial zone for a flow parallel to the edge of the filler as a function of the distance to the outer edge of the interface zone x,  $D_{Imax}$  is the maximum of the diffusivity in the interfacial zone, and  $D_M$  is the diffusivity in the matrix. x is a control variable that represents different relative distances from the outer edge of the interface at which the diffusivity is calculated. This is necessary because in ABAQUS the material data can only be provided as a table and not as a function. In order to generate the desired course of the diffusion rate change, a prefactor  $\Psi$  was introduced. This type of function was chosen because it allows for a large gradient without creating jumps

Algorithm 2: Pseudocode example for an ABAQUS UFIELD subroutine.				
SUBROUTINE UFIELD				
/* The data of the node for wich this subroutine is called. */				
Input: nodeData;				
/* Create a global list, which can be called from every				
subroutine call, in which the data of all the filler				
particles is stored. */				
<b>Create Global List:</b> fillerPartikelList = EMPTY;				
/* Create a variable to hold numerical data. */				
<b>Create Variable:</b> minDistanceToFillerParticle = INFINITY;				
/* If the filler particle list is empty, load the necessary				
data from a csv file. This is only necessary for the first				
time this subroutine is called. */				
if fillerPaticelList == Empty then $f(t) = f(t) = f(t)$				
<b>Ioad</b> fillerParticleData <b>Into</b> fillerPartikelList;				
/* Search for the distance of the closest filler particle to				
the node. */				
for fillerParticle in fillerPartikelList do				
<b>calculate</b> distance <b>from</b> node.position <b>to</b> fillerParticle;				
If distance < minDistanceToFillerParticle then				
/* Calculate the relative position of the node in relation to				
the thickness of the interface zone. */				
calculate relativePositionValue;				
/* Save the relative position as a field variable. */				
<b>Save</b> relativePositionValue in FIELD;				

in the curve. Various experiments were then carried out to find a  $\Psi$  for which the curve has no jumps at the beginning, a large rate of change while still maintaining the bleeding effect, which was explained in Section 3.1. Such curves are shown in Fig. 3.6. After evaluating these experiments, it was decided to use a  $\Psi$  of 20 for the simulations in this thesis.

Furthermore, for future projects, the adhesion coefficient mentioned in the introduction can be represented by the two variables  $D_{Imax}$  and  $\Psi$ .



**Fig. 3.6** Curve shapes of diffusion rates in the interface zone at different prefactors  $\Psi$ . A value of 0 on the x-axis represents the position at the outer edge of the interface zone and a value of 1 represents the position at the outer edge of the filler particle.

## 3.3 Simulation Workflow

#### 3.3.1 Overview

Since the used simulation models differ only in geometry and the resulting relevant code differences have already been explained in Section 3.1, the proposed simulation workflow is generally valid for this work. First of all, each simulation has the following basic configuration:

- The inlet of the membrane is at the top edge. There the concentration of the permeate is constant at a value of  $1 \text{ mol} \cdot \text{m}^{-2}$ .
- The outlet of the membrane is at the bottom edge. There the concentration of the permeate is constant at a value of  $0 \text{ mol} \cdot \text{m}^{-2}$ .
- The initial concentration of the permeate in the membrane is set to  $0 \text{ mol} \cdot \text{m}^{-2}$ .
- The left and right edges of the membrane are connected with a periodic boundary condition.
- The filler particle is always implemented as a hole in the membrane. The permeate can therefore not penetrate the particle.

• Every filler particle has an interface zone around it, as described in Section 3.1.

The general workflow of the simulations used in this work consists of multiple steps, which are controlled by one central script. First of all, the following simulation parameters are defined in this central script:

- Filler particle thickness, *b*
- Filler particle aspect ratio,  $\alpha$
- Filler particle area ratio,  $\phi_f$
- Number of filler particles in the membrane
- Mesh seeding sizes of the membrane, mesh transition zone and the interface zone
- Inlet and outlet concentration of the diffusion medium,  $C_1$  and  $C_2$ , respectively
- Diffusion coefficients for both materials,  $D_M$  for the matrix material,  $D_{IP}$  and  $D_{IN}$  for the interface material
- Solubility coefficient for both materials,  $S_M$  for the matrix and  $S_I$  for the interface

Based on these simulation parameters, the positions of all fillers in the membrane are then calculated, inserted into a list and saved together with the simulation parameters in CSV files. These files are then automatically read in during the simulation steps in which they are required. Next, the following steps are performed:

- 1. Generation of the ABAQUS input file based on the parameters from the CSV file.
- 2. Running the ABAQUS interface simulation, with the subroutines.
- 3. Extracting data out of the ABAQUS ODB file.
- 4. Running the ABAQUS standard simulation.

These steps are explained in detail in the following subsections. At the end of a simulation series, the result files were read in again with the help of a subsequent script, the data contained therein was transformed into the desired form and plotted in various plots. The actual code used in this Python script is provided in the appendix in Section A.2.4.

#### **3.3.2** Creation of Input File

First, the ABAQUS input file is created on the basis of the simulation parameters defined in the input. For this purpose ABAQUS is automated by a Python script which runs the following sequence:

- 1. Geometry creation: The membrane is drawn as a rectangle. The filler particles are simultaneously cut into this rectangle as holes. The area around these holes is then partitioned twice. Once at a very small distance from the perimeter of the hole and once at ten times that distance.
- 2. Material creation: The interface and membrane materials are created and assigned to the corresponding areas of the model.
- 3. Mesh creation: The membrane mesh is kept coarse, while the edges of the filler particles are seeded with a finer mesh size. The outer edges of the interface zone are also seeded with the same fine mesh size. The outer edges of the mesh transition zone are seeded with an intermediate mesh size between the size of the membrane and the interface zone mesh. This results in a mesh configuration as shown in Fig. 3.7. This mesh configuration is necessary for the above mentioned subroutines to work with sufficient accuracy.
- 4. Field creation: The ABAQUS field for the UFIELD Subroutine is created and assigned to the interface zone.
- 5. Set creation: The necessary ABAQUS node sets on the left, right and bottom edge are created. The sets on the left and right edge are necessary for the periodic boundary condition which connects these two edges. The set on the bottom edge is necessary for collecting the output data.
- 6. Time step creation: A time step is the framework for the boundary conditions. That means in this step the inlet and outlet boundary conditions as well as the periodic boundary are defined.

The actual code used in the Python script is provided in the appendix in Section A.2.1.

#### 3.3.3 ABAQUS Interface Simulation

After the ABAQUS input file has been created, the simulation is started by a batch file. This batch file is needed to load the input file into the simulations and to connect them with the
## **3** Simulation



Fig. 3.7 Mesh of membranes with (a) a circular filler particle and (b) a rectangular filler particle.

subroutines. When the simulation is finished, ABAQUS is started again automated by a Python script and the ODB file of the simulation is opened. Afterwards, the script extracts all relevant data and saves them to a CSV file. This Python script is provided in the appendix in Section A.2.2.

## 3.3.4 ABAQUS Standard Simulation

For the standard simulation, another simulation is performed with the previously defined simulation parameters, but without an interface zone material in the model. However, the mesh is generated in the same way as for the simulation with an interface zone to avoid a significant difference in the results. Since no subroutines are required for this simulation, all three steps, i.e. model creation, simulation and data extraction, can be automated using a single Python script. This simulation is necessary to compare the results of the interface simulations. The actual code used in the Python script is provided in the appendix in Section A.2.3.

## Chapter 4

# Results

This chapter presents the results of the simulations performed in the course of this thesis. Two types of plots have been chosen to highlight different aspects of the results. First, the raw results of the different simulations are discussed to show the influence of the different parameters on the resulting outflow over the width of the membrane. Then, the outflow values are integrated over the individual membranes to obtain the total outflow for a given membrane. These integrated values are better suited for comparison with one another. Furthermore, some of the results obtained from ABAQUS are shown directly as examples.

## 4.1 **Raw Results - One Filler Particle**

In Fig. 4.1 the permeate concentration is shown for the simulations with only one central filler particle in the membrane, as provided by ABAQUS. Comparing Fig. 4.1a with Fig. 4.1b and Fig. 4.1c with Fig. 4.1d it can be observed that the implementation of an interface zone leads to a better distribution of the permeate around the filler. It can be clearly seen that the interface zone increases the flow of the medium through the membrane.

In Fig. 4.2, the raw results of simulations with a circular and a rectangular filler particle are shown. These simulations were carried out with the following parameters:

- Diffusivity of the matrix material,  $D_M = 1.0 \text{ m}^2 \cdot \text{s}^{-1}$
- Diffusivity of the interface material parallel to the filler particle edges,  $D_{IP} = 10.0 \text{ m}^2 \cdot \text{s}^{-1}$
- Diffusivity of the interface material normal to the filler particle edges,  $D_{IN} = 1.0 \text{ m}^2 \cdot \text{s}^{-1}$



**Fig. 4.1** Concentration results of the simulations of a membrane with (a) a circular filler in the center and an interface zone, (b) a circular filler in the center and no interface zone, (c) a rectangular filler in the center and an interface zone and (d) a rectangular filler in the center and no interface zone.

• Interface thickness,  $w_I = 0.5 \,\mu\text{m}$ 

The purpose of these simulations was to get an overview of the influence of  $\Psi$  on the simulation results. This was necessary as the value of  $\Psi$  was only assigned based on the results in Fig. 3.6 and a validation was required. As can be seen from these figures, a  $\Psi$  of 50 or more results in a complete overlap of the curves. The reason for this is that only the diffusivity of the mesh nodes closest to the filler particle is affected. However, since it was desired to emulate a bleeding effect as described in Section 3.1, it was decided to leave the value of  $\Psi$  at 20. These figures also show the effect of the central filler particle on the outflow from the membrane. In both cases (circle and rectangle), there is a significant drop in the outflow at the center of the membrane. However, the width of this drop in the graph is greater for the rectangular filler particle. This is due to the fact that the dimensions of the round and



rectangular fillers were chosen so that the areas are equal, and since the rectangular particle is not square, the result is a particle that is wider perpendicular to the direction of diffusion.

**Fig. 4.2** (a) Outflow F over the width of a membrane with one circular filler in the middle, (b) outflow over the width of a membrane with one rectangular filler in the middle.

## 4.2 Raw Results - Multiple Filler Particles

In Fig. 4.3, the concentration results are shown for the simulations with multiple rectangular filler particles in the membrane, as provided by ABAQUS. Comparing Fig. 4.3a with Fig. 4.3b shows that the size of the filler particles is too small compared to the membrane to produce any visible differences. This proves the previous point that it is necessary to perform preliminary simulations on simpler systems to gather information for the more complex simulation.

Fig. 4.4 shows the raw results of simulations with multiple filler particles in the membrane. These simulations were carried out with the following parameters:

• 
$$\phi_f = 0.01$$



**Fig. 4.3** Concentration results of simulating a membrane with (a) multiple rectangular filler particles and an interface zone around each filler, (b) multiple rectangular filler particles without the interface zone.

- $w_I = 0.25 \,\mu m$
- $D_M = 1.0 \,\mathrm{m}^2 \cdot \mathrm{s}^{-1}$
- $D_{IP} = 10.0 \,\mathrm{m}^2 \cdot \mathrm{s}^{-1}$
- $D_{IN} = 1.0 \,\mathrm{m}^2 \cdot \mathrm{s}^{-1}$
- Interface Prefactor ( $\Psi$ ) = 20.0

The results shown here are an excerpt from the final results of this work, which are shown and described in Fig. 4.6. The purpose of this presentation is to show how the outflow of a



Fig. 4.4 Outflows across the width of various membranes with multiple rectangular fillers. The membranes differ by the  $\alpha$  values of the filler particles, as shown in the legends of the plots.

membrane with multiple filler particles behaves. The width of the membrane depends on the number of filler particles, the aspect ratio  $\alpha$  of the particles, and  $\phi_f$  of the particles in the membrane. Therefore, as the aspect ratio increases, the width of the membrane also increases, as shown in Fig. 4.6. The thickness of the filler particles, i.e. their expansion in the diffusion direction, is constant, so  $\alpha$  only affects the width of the filler particles, i.e. their expansion perpendicular to the diffusion direction. This is the reason why the valleys in the curves in Fig. 4.6 become wider as  $\alpha$  increases. Furthermore, this behavior of the filler particles also results in a stronger barrier effect, which also reduces the total amount of outflow Q as  $\alpha$  increases.

## 4.3 Accumulated Permeate - One Filler Particle

In Fig. 4.5 the processed results of several simulations with either a round or a rectangular filler particle are shown. The results were processed by integrating the outflow values over the width of each membrane and accumulating the outflow values over the duration of the simulation to obtain the total outflow over time  $Q_{total}$  from the membrane. These results show the effect of  $\Psi$  and  $D_{IP}$  on the amount of total outflow from the membrane.

When comparing the two subfigures in Fig. 4.5, it is noticeable that the integrated outflow values of the membrane with a round filler particle are higher than those of the membrane

with a rectangular filler particle, although both particle types have the same area. This difference is due to the fact that the rectangular particle is wider than the round particle. To achieve a higher barrier effect for the same filler area fraction, particles with a large width but small thickness should be used. In the case of the definition used in this work, this means that a high aspect ratio results in better barrier properties for the same filler area fraction. Under this condition, a significant barrier effect can be achieved, even with a low particle filler area fraction. However, highly aspected filler particles have the disadvantage that the orientation of the filler particles has a major influence on the barrier effect.

The resulting value of a simulation without interface zone around the filler particle is shown in both subfigures of Fig. 4.5 as a red line. If  $D_{IP}$  is set to 1, the resulting cumulative outflow values lie on the red line. If  $D_{IP}$  is greater than 1, the result is above the red line, indicating that more outflow has occurred. If  $D_{IP}$  is less than 1, the result is below the red line, indicating less outflow.



**Fig. 4.5** (a) Integrated outflow of membranes with one circular filler in the middle, (b) integrated outflows of membranes with one rectangular filler in the middle.

## 4.4 Accumulated Permeate - Multiple Filler Particles

One aim of this work was to compare the following three types of data and assess the influence of an interface zone on the diffusivity of reinforced membranes:

- Analytically obtained data from the Nielsen model described in detail in Section 2.5
- FEM simulated result data with interface zones, which was described in detail in Section 3.3.3
- FEM simulated result data without interface zones, which was described in detail in Section 3.3.4

This comparison is shown in Fig. 4.6, where the analytical data are plotted as continuous lines, the FEM data without interface zone are labeled "o" and the FEM data with interface zone are labeled "x". As described in Section 2.5, the filler particles in the Nielsen model are arranged in a regular pattern. This creates channels in which unhindered diffusion flow is possible. As  $\phi_f$  and  $\alpha$  increase, the width of these channels relative to the membrane width decreases. Above a certain  $\phi_f$  value, this ratio remains more or less constant. The  $\phi_f$  value at which this plateau occurs, depends on  $\alpha$ . The higher the  $\alpha$  value, the earlier the plateau occurs. This structure of the membrane causes the outflow to decrease with increasing  $\phi_f$  as well as with increasing  $\alpha$ . This result was expected because increasing the number of filler particles or increasing the size of the filler particles should result in an increased barrier effect.

Since the differences between the FEM data and the analytical data are difficult to see in Fig. 4.6, the absolute deviations of the FEM data from the analytical data were calculated and are shown in Fig. 4.7. The standard FEM model agrees quite well with the analytical model for filler area fractions between 0.06 and 0.09. When  $\phi_f$  is below this range, the results of the FEM model are higher than those of the analytical model. When  $\phi_f$  is above this range, the results are lower than those of the analytical model. The results of the FEM model with interface zone are consistently higher than those of the standard FEM model, since the interfacial zone facilitates diffusion flow through the membrane. The advanced FEM model agrees quite well with the analytical model for filler area fractions between 0.06 and 0.15.



Fig. 4.6 Accumulated outflows of membranes with multiple rectangular fillers.



Fig. 4.7 Absolute error of the total outflow from the membranes in FEM models in respect to the analytically obtained data.

# Chapter 5

# Conclusions

The aim of this thesis was to numerically simulate the diffusion of hydrogen through different particle filled membranes with the introduction of a filler-matrix interface. The difference in the membranes was based on the area fraction of the particles in the membrane and the width of the filler particles. ABAQUS was chosen as the numerical solver for this work because of the software's feature set and in-house knowledge of it. The extended Nielsen model was selected as the basis for the simulation model because it has the advantage that it can be used in both analytical and numerical simulations, making the results directly comparable.

Implementing the interface zone was difficult for two reasons. Firstly, the implementation of an interface zone in ABAQUS was not easy. Instead of the ABAQUS GUI, Fortran subroutines had to be used, which had to be learned first. Furthermore, a lot of research had to be done to find the right subroutines for the problem at hand, as there are over 100 different subroutines avaiable in ABAQUS. Each of these subroutines has its own usecase, comes with different methods and functions and has access to different systems of ABAQUS. Secondly, the implementation of the Python modules and the subroutines had to be improved, as a first, non-optimized attempt showed very poor performance due to too many file calls for each calculation step. All calculations necessary for the subroutines to function were carried out using Python scripts and the results were saved on the hard disk. These result files were then opened several times by the Fortran subroutines in each simulation step.

In its final implementation, the interface zone was integrated into the model at a mesh node and integration point level. Three subroutines were used for this setup. The first subroutine had the sole purpose of calculating the orientation of each node in the interface zone, based on its position in this zone, relative to the global coordinate system. The second subroutine was used to calculate the position of each node in the interface zone relative to the thickness

## **5** Conclusions

of the interface zone. Therefore, if the node is exactly at the edge of the filler particle, the relative position value is equal to one, and if the node is at the outer edge of the interface zone, the relative position value is equal to zero. The results of these calculations are stored in a scalar field attached to all interface mesh nodes. The third subroutine is necessary because the relative position values stored in the scalar field cannot be used directly by ABAQUS to adjust the material data. This third subroutine takes the scalar field values at the mesh node positions and interpolates them to the integration points of the corresponding mesh elements. With this sophisticated system in place, it was possible to create the desired behavior of the interface zone and perform the necessary simulations for this work.

The results of this work not only show that such a FEM model with interface zone is possible, but also provide plausible data with respect to the standard FEM and the analytical model. However, both numerical models show larger deviations from the analytical model for membranes with low filler area fraction and small filler particle widths. The analytical model predicts better barrier properties for such filled membranes than the numerical simulations. Due to the lack of data from real experiments, it was not possible to quantitatively compare the two numerical models. After completing the work, it can be said that ABAQUS has a steep learning curve, but the simulation results and the adaptability of the simulation models with the subroutines are convincing.

# **Chapter 6**

# Outlook

The next step to improve this model should be to optimize the code to achieve better simulation runtime and stability. Before extending the simulation setup further, consideration should be given to comparing the results not only with analytical data, but also with real experimental results. These comparisons are essential for the evaluation of the model. Below are some suggestions about additions to the model that could lead to an improvement in accuracy.

For example, the filler particles could be placed randomly. This method of membrane generation would allow for more naturally built models. Another way to generate more realistic data would be to develop a three-dimensional system. This setup would make it easier to replicate real world experiments. Since the Nielsen model, which was heavily referenced in this thesis, only works in two dimensions, either a new model for the structure of the membrane would have to be developed, the filler particles would have to be placed randomly or the model could be based on a CT scan of a real reinforced polymer membrane.

For a different simulation model in this project, an attempt was made to implement a randombased distribution of filler particles in a two-dimensional model. The implementation was successful, but required a lot of optimization to achieve a reasonable runtime of the membrane generation. It can be concluded that it would be even more difficult to achieve a reasonable implementation in a three-dimensional model. [26]

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Aspect ratio of the filler particles.	(dimensionless)
Change of the concentration in the mem-	$\left[\frac{\partial C}{\partial t}\right] = \operatorname{mol} \cdot \operatorname{m}^{-2} \cdot \operatorname{s}^{-1}$
brane over time. (2D)	
Pressure gradient along the thickness of	$\left[\frac{\partial p}{\partial x}\right] = \operatorname{Pa} \cdot \mathrm{m}^{-1}$
the membrane. (2D)	
Dimensionless concentration of the gas in	(dimensionless)
the membrane.	
Dimensionless time.	(dimensionless)
Dimensionless distance.	(dimensionless)
Number of filler particles in the mem-	(dimensionless)
brane.	
Mobility of the particles.	$[\mu] = s \cdot kg^{-1}$
Gradient of the concentration along the	$[\nabla C] = \operatorname{mol} \cdot \mathrm{m}^{-3}$
thickness of the membrane. (2D)	
Area fraction of the filler particles in the	(dimensionless)
membrane.	
Diffusion prefactor, necessary to calcu-	(dimensionless)
late the interface diffusion coefficent in	
respect to the relative position in the inter-	
face zone.	
Tortuousity factor, describes the longer	(dimensionless)
tortuous path of the diffusing medium due	
to the impermeable filler particles.	
Area of the filler particle in an unit cell.	$[A_f] = m^2$
Area of a unit cell in the membrane.	$[A_u] = \mathrm{m}^2$
Thickness of the filler particle.	[ <i>b</i> ] = m
	Aspect ratio of the filler particles. Change of the concentration in the mem- brane over time. (2D) Pressure gradient along the thickness of the membrane. (2D) Dimensionless concentration of the gas in the membrane. Dimensionless time. Dimensionless distance. Number of filler particles in the mem- brane. Mobility of the particles. Gradient of the concentration along the thickness of the membrane. (2D) Area fraction of the filler particles in the membrane. Diffusion prefactor, necessary to calcu- late the interface diffusion coefficent in respect to the relative position in the inter- face zone. Tortuousity factor, describes the longer tortuous path of the diffusing medium due to the impermeable filler particles. Area of the filler particle in an unit cell. Area of a unit cell in the membrane. Thickness of the filler particle.

$b_H$	Describes the gradient of the linear seg- ment at the start of the Langmuir model.	$[b_H] = \mathrm{Pa}^{-1}$
С	Concentration of the gas in the membrane. (2D)	$[C] = \operatorname{mol} \cdot \operatorname{m}^{-2}$
$C'_H$	Membrane saturation constant in the Langmuir model. (2D)	$[C'_H] = \mathrm{mol} \cdot \mathrm{m}^{-2}$
$C_1$	Concentration of the gas at the upstream border of the membrane. (2D)	$[C_1] = \operatorname{mol} \cdot \operatorname{m}^{-2}$
<i>C</i> <sub>2</sub>	Concentration of the gas at the down- stream border of the membrane. (2D)	$[C_2] = \operatorname{mol} \cdot \operatorname{m}^{-2}$
D	Diffusion coefficient of the membrane. (2D)	$[D] = \mathbf{m}^2 \cdot \mathbf{s}^{-1}$
d	Distance between two rows of filler particles.	[ <i>d</i> ] = m
$D_0$	Diffusion coefficient of the unfilled mem- brane. (2D)	$[D_0] = \mathbf{m}^2 \cdot \mathbf{s}^{-1}$
Deff	Effective diffusion coefficient of the filled membrane. (2D)	$[D_{eff}] = \mathbf{m}^2 \cdot \mathbf{s}^{-1}$
D <sub>Imax</sub>	Maximum diffusion coefficient in the in- terface zone. (2D)	$[D_{Imax}] = \mathrm{m}^2 \cdot \mathrm{s}^{-1}$
$D_{IN}$	Diffusion coefficient in the interface zone, if flow is normal to the filler particle edge. (2D)	$[D_{IN}] = \mathrm{m}^2 \cdot \mathrm{s}^{-1}$
D <sub>IP</sub>	Diffusion coefficient in the interface zone, if flow is parallel to the filler particle edge. (2D)	$[D_{IP}] = \mathrm{m}^2 \cdot \mathrm{s}^{-1}$
$D_M$	Diffusion coefficient of the membrane out- side of the interface zone. (2D)	$[D_M] = \mathrm{m}^2 \cdot \mathrm{s}^{-1}$
F	Diffusive flow density through the mem- brane. (2D)	$[F] = \operatorname{mol} \cdot \operatorname{s}^{-1} \cdot \operatorname{m}^{-1}$
$F_0$	Flow density through the unfilled mem- brane. (2D)	$[F_0] = \operatorname{mol} \cdot \mathrm{s}^{-1} \cdot \mathrm{m}^{-1}$
<i>F<sub>ch</sub></i>	Unhindered flow density through the chan- nels of the membrane. (2D)	$[F_{ch}] = \operatorname{mol} \cdot \operatorname{s}^{-1} \cdot \operatorname{m}^{-1}$
f <sub>ch</sub>	Ratio of the flow through the unobstructed channel in respect to the total flow.	(dimensionless)

Feff	Effective flow through the filled mem- brane (2D)	$[F_{eff}] = \operatorname{mol} \cdot \mathrm{s}^{-1} \cdot \mathrm{m}^{-1}$
$F_t$	Tortuous flow density through the chan- nels of the membrane (2D)	$[F_t] = \operatorname{mol} \cdot \operatorname{s}^{-1} \cdot \operatorname{m}^{-1}$
$f_t$	Ratio of the tortuous flow in respect to the total flow.	(dimensionless)
kh	Boltzmann constant.	$[k_h] = \mathbf{J} \cdot \mathbf{K}^{-1}$
Ľ	Membrane width.	[L] = m
L <sub>eff</sub>	Effective length of the tortuous path of the permeate through the filled membrane.	$[L_{eff}] = \mathbf{m}$
$L_t$	Length of the tortuous path through the unit cell.	$[L_t] = \mathbf{m}$
Р	Permeation coefficient of the membrane.	$[P] = \operatorname{mol} \cdot \operatorname{s}^{-1} \cdot \operatorname{Pa}^{-1}$
р	Partial pressure of the gas in the mem- brane.	[p] = Pa
$p_1$	Pressure of the gas at the upstream bound- ary of the membrane.	$[p_1] = Pa$
<i>p</i> <sub>2</sub>	Pressure of the gas at the downstream boundary of the membrane.	$[p_2] = Pa$
Q	Cumulative diffusive flow density through the membrane. (2D)	$[Q] = \operatorname{mol} \cdot \operatorname{m}^{-1}$
$Q_{total}$	Cumulative diffusive flow through the membrane.	$[Q_{total}] = mol$
S	Solubility coefficient of the membrane. (2D)	$[S] = \operatorname{mol} \cdot \operatorname{m}^{-2} \cdot \operatorname{Pa}^{-1}$
S	Slit shape, distance between two filler par- ticles in the same row.	[s] = m
$S_0$	Solubility coefficient of the unfilled mem- brane. (2D)	$[S_0] = \operatorname{mol} \cdot \operatorname{m}^{-2} \cdot \operatorname{Pa}^{-1}$
$S_{eff}$	Effective solubility coefficient of the filled membrane. (2D)	$[S_{eff}] = \mathrm{mol} \cdot \mathrm{m}^{-2} \cdot \mathrm{Pa}^{-1}$
$S_I$	Solubility coefficient in the interface zone.	$[S_I] = \operatorname{mol} \cdot \operatorname{m}^{-2} \cdot \operatorname{Pa}^{-1}$
$S_M$	Solubility coefficient in the undisturbed membrane. (2D)	$[S_M] = \operatorname{mol} \cdot \operatorname{m}^{-2} \cdot \operatorname{Pa}^{-1}$
Т	Temperature.	[T] = K

## 10 List of symbols

$T_c$	Critical temperature of a gas.	$[T_c] = K$
$T_g$	Glass transition temperature of a polymer.	$[T_g] = K$
W	Width of the filler particle.	[w] = m
WI	Thickness of the interface zone.	$[w_I] = m$
$W_{ch}$	Width of a channel in a filler particle ar-	$[w_{ch}] = m$
	rangement.	

# Appendix A

# Codes

## A.1 Fortran - ABAQUS Subroutines

## A.1.1 Subroutine ORIENT

```
SUBROUTINE ORIENT(T, NOEL, NPT, LAYER, KSPT, COORDS, BASIS,
     * ORNAME, NNODES, CNODES, JNNUM)
2
3 C
        INCLUDE 'ABA_PARAM. INC'
4
 С
5
        CHARACTER*80 ORNAME
6
7 C
        DIMENSION T(3,3), COORDS(3), BASIS(3,3), CNODES(3, NNODES)
8
        DIMENSION JNNUM(NNODES)
9
10
        user coding to define T
11
12
        RETURN
13
        END
14
```

## A.1.2 Subroutine UFIELD

```
SUBROUTINE UFIELD (FIELD, KFIELD, NSECPT, KSTEP, KINC, TIME, NODE,
     * COORDS, TEMP, DTEMP, NFIELD)
 С
3
       INCLUDE 'ABA_PARAM. INC'
4
5 C
       DIMENSION FIELD(NSECPT, NFIELD), TIME(2), COORDS(3),
6
      1 TEMP(NSECPT), DTEMP(NSECPT)
7
 C
8
9
10
       user coding to define FIELD
11
13
       RETURN
14
       END
15
```

## A.1.3 Subroutine USDFLD

```
SUBROUTINE USDFLD(FIELD, STATEV, PNEWDT, DIRECT, T, CELENT,
     * TIME, DTIME, CMNAME, ORNAME, NFIELD, NSTATV, NOEL, NPT, LAYER,
     * KSPT, KSTEP, KINC, NDI, NSHR, COORD, JMAC, JMATYP, MATLAYO, LACCFLA)
3
 С
4
        INCLUDE 'ABA PARAM. INC'
5
 С
6
        CHARACTER*80 CMNAME, ORNAME
7
        CHARACTER*3 FLGRAY(15)
8
        DIMENSION FIELD(NFIELD), STATEV(NSTATV), DIRECT(3,3),
9
     * T(3,3),TIME(2)
10
        DIMENSION ARRAY(15), JARRAY(15), JMAC(*), JMATYP(*), COORD(*)
11
13
         user coding to define FIELD and, if necessary, STATEV and PNEWDT
14
15
16
        RETURN
17
        END
18
```

## A.2 Python - ABAQUS Control Scripts and General Utility Scripts

## A.2.1 ABAQUS Input File Creation

```
# Imports
 from abaqus import *
2
3 from abaqusConstants import *
 from caeModules import *
6 import csv
  import math
  import numpy as np
  import re
9
10 import sys
12
13 def log(msg):
      # Helper-function to print logs to the python cosole, the abaqus
14
     console and save them in a log file
      print(str(msg))
15
      print >> sys.__stdout__, str(msg)
16
      with open("../Logs/log.txt", "a") as log_file:
          log_file . write ( str (msg) )
18
          log_file . write("\n")
19
20
22 # Variables
23 case = None
  with open("../DATA_FILES/actual_case.txt") as file:
24
      reader = csv.reader(file, delimiter=",")
25
      for row in reader:
26
          if row:
27
               case = row[0].replace("\\", "/").replace("DATA_FILES/", "")
28
29
30
31 case_path = ".../DATA_FILES/" + case
32
33 geometric_params = []
34 with open(case_path + "/Geometric_Data.csv") as file:
      reader = csv.reader(file, delimiter=",")
35
      for row in reader:
36
```

```
for row_elem in row:
37
               str_floats = re.findall(r"[-+]?\d*\.\d+|\d+", row_elem)
38
39
               tmp = []
               for str elem in str floats:
40
                   tmp.append(float(str_elem))
41
               geometric_params.append(tmp)
42
43
44 name_params = None
  with open(case_path + "/Output_Data.csv") as file:
45
      reader = csv.DictReader(file, delimiter=",")
46
      for row in reader:
47
           name_params = row
48
49
50 filler_radius = geometric_params[0][0]
si interface_thickness = geometric_params[1][0]
52 membrane_thickness = geometric_params [2][0]
53 membrane_width = geometric_params[3][0]
54
solubility = geometric_params[4][0]
56 diffusivity_matrix = geometric_params[5][0]
57 diffusivity_interface = geometric_params[6][0]
58 interface_prefactor = geometric_params[7][0]
59
60 matrix_mesh_size = geometric_params[8][0]
 interface_mesh_size = geometric_params[9][0]
61
62
 filler_center = [geometric_params[10][0], geometric_params[11][0]]
63
64
  interface_distances = [
65
      0.05,
66
      0.1,
67
      0.15,
68
      0.2,
69
      0.25,
70
      0.3,
71
      0.35,
72
      0.4,
73
      0.45,
74
      0.5,
75
      0.55,
76
      0.6,
77
      0.65,
78
      0.7,
79
      0.75,
80
```

```
0.85,
81
      0.9,
82
83
      0.95,
      0.98,
84
      0.99,
85
  1
86
  diffusivity = [(1.0, diffusivity_matrix, 0.0, 0.0, 0.0)]
87
88
  for dist in interface_distances:
89
      diff_value = diffusivity_interface * math.exp(interface_prefactor *
90
      (dist - 1)) + diffusivity_matrix * (
           1 - math.exp(interface_prefactor * (dist - 1))
91
      )
92
      diffusivity.append((1.0, diff_value, 0.0, 0.0, dist))
93
94
  diffusivity.append((1.0, diffusivity_interface, 0.0, 0.0, 1.0))
95
  diffusivity = tuple(diffusivity)
96
97
  # _____
98
  # Abaqus Setup
99
  # _____
100
101
102 Mdb()
my_session = session
104 \text{ my_mdb} = \text{mdb}
105
  my_mdb.models.changeKey(fromName="Model-1", toName=name_params["model"])
106
  my_model = my_mdb.models[name_params["model"]]
107
108
  # _____
109
110 # Part Creation
  # _____
sketch = my_model. ConstrainedSketch(name="__profile__", sheetSize=200.0)
sketch.rectangle(point1=(0, 0), point2=(membrane_width,
      membrane_thickness))
115 p1 = (filler_center[0] + filler_radius, filler_center[1])
116 sketch. CircleByCenterPerimeter(center=filler_center, point1=p1)
my_part = my_model. Part (name=name_params ["part"], dimensionality=
     TWO_D_PLANAR, type=DEFORMABLE_BODY)
118 my_part. BaseShell(sketch=sketch)
119
120 sketch = my_model. ConstrainedSketch(name="__profile__", sheetSize=200.0)
```

```
121 p2 = (filler_center[0] + filler_radius + interface_thickness ,
      filler center [1])
  sketch.CircleByCenterPerimeter(center=filler_center, point1=p2)
  my_part.PartitionFaceBySketch(faces=my_part.faces, sketch=sketch)
124
  #
    _____
125
  # Material Setup
126
  #
128
  my_material = my_model.Material(name=name_params["material_matrix"])
129
130 my_material. Diffusivity (law=GENERAL, table=((diffusivity_matrix,),))
  my_material. Solubility (table = ((solubility ,) ,))
  my_material = my_model. Material (name=name_params ["material_interface"])
  my_material.Diffusivity(type=ORTHOTROPIC, law=GENERAL, dependencies=1,
134
      table=diffusivity )
  my_material. Solubility (table = (( solubility ,) ,) )
135
  my_material.UserDefinedField()
136
  #
   _____
138
  # Section Setup
139
  #
140
    _____
141
  my_model.HomogeneousSolidSection(
142
      name=name_params["section_matrix"], material=name_params["
143
      material_matrix"], thickness=None
144
  )
  my_model.HomogeneousSolidSection(
145
      name=name_params["section_interface"], material=name_params["
146
      material_interface"], thickness=None
147
  )
148
  region_matrix = regionToolset.Region(
149
      faces=my_part.faces.findAt(
150
           ((filler_center[0] + filler_radius + interface_thickness + 1e-3,
151
       filler_center [1], 0),),
      )
  )
153
  region_interface = regionToolset.Region(
154
      faces=my_part.faces.findAt(
155
           ((filler_center[0] + filler_radius + interface_thickness - 1e-3,
156
       filler_center [1], 0),),
      )
157
158)
```

## A Codes

```
my_part.SectionAssignment(
159
       region=region_matrix,
160
161
       sectionName=name_params["section_matrix"],
       offset = 0.0,
162
       offsetType=MIDDLE_SURFACE,
163
       offsetField="",
164
       thicknessAssignment=FROM_SECTION,
165
  )
166
  my_part.SectionAssignment(
       region=region_interface,
168
       sectionName=name_params["section_interface"],
169
       offset = 0.0,
170
       offsetType=MIDDLE_SURFACE,
       offsetField="",
       thicknessAssignment=FROM_SECTION,
174
  )
175
  #
176
  # Mesh Setup
177
   _____
  #
178
179
  elemType1 = mesh.ElemType(elemCode=DC2D4, elemLibrary=STANDARD)
180
  elemType2 = mesh.ElemType(elemCode=DC2D3, elemLibrary=STANDARD)
181
182
  my_part.setMeshControls(regions=my_part.faces, elemShape=QUAD, algorithm
183
      =ADVANCING_FRONT)
  my_part.setElementType(
184
       regions=regionToolset.Region(faces=my_part.faces),
185
       elemTypes =(elemType1, elemType2),
186
  )
187
  my_part.seedPart(size=matrix_mesh_size, deviationFactor=0.1,
188
      minSizeFactor = 0.1)
189
  interface_edges = my_part.edges.getByBoundingBox(
190
      xMin=(filler_center[0] - filler_radius - interface_thickness - 1e-3)
191
      xMax=(filler_center[0] + filler_radius + interface_thickness + 1e-3)
192
      yMin=(filler_center[1] - filler_radius - interface_thickness - 1e-3)
193
      yMax=(filler_center[1] + filler_radius + interface_thickness + 1e-3)
194
195
  )
196 my_part.Set(edges=interface_edges, name="SeedByEdge")
```

```
my_part.seedEdgeBySize(edges=interface_edges, size=interface_mesh_size)
197
  my_part.generateMesh()
198
199
  #
    _____
200
    Oriantation Setup
  #
201
  #
    _____
202
203
  my_part. MaterialOrientation (region=region_interface, orientationType=
204
      USER)
205
  # _____
206
  # Assembly Setup
207
    _____
  #
208
209
  my_assembly = my_model.rootAssembly
210
  my_assembly.DatumCsysByDefault(CARTESIAN)
  my_instance = my_assembly.Instance(name=name_params["part"] + "-1", part
213
      =my_part, dependent=ON)
214
  # _____
215
  # Set Setup
216
    _____
  #
217
218
  all_nodes = my_instance.nodes
219
  top_nodes = my_instance.nodes.getByBoundingBox(
220
       xMin = -1e - 3,
221
       yMin=membrane_thickness - 1e-3,
       zMin = 0.0,
       xMax=membrane_width + 1e-3,
224
      yMax=membrane_thickness + 1e-3,
225
      zMax = 0.0,
226
  )
  bot_nodes = my_instance.nodes.getByBoundingBox(
228
       xMin=-1e-3, yMin=-1e-3, zMin=0.0, xMax=membrane_width + 1e-3, yMax=1
229
      e - 3, zMax = 0.0
  )
230
  interface_nodes = my_instance.nodes.getByBoundingCylinder(
       center1 = (filler_center[0], filler_center[1], 0 - 1e-3),
233
       center2 = (filler_center[0], filler_center[1], 0 + 1e-3),
234
       radius=filler_radius + interface_thickness + 1e-3,
235
236)
237
```

```
interface_elements = my_instance.elements.getByBoundingCylinder(
238
       center1 = (filler_center[0], filler_center[1], 0 - 1e-3),
239
       center2 = (filler_center[0], filler_center[1], 0 + 1e-3),
240
       radius=filler_radius + interface_thickness + 1e-3,
241
  )
242
243
  all_nodes_set = my_assembly.Set(nodes=all_nodes, name=name_params["
244
      set_All"])
  interface_nodes_set = my_assembly.Set(nodes=interface_nodes, name=
245
      name_params["set_Interface_Nodes"])
  interface_element_set = my_assembly.Set(nodes=interface_elements, name=
246
      name_params["set_Interface_Elements"])
  inlet_set = my_assembly.Set(nodes=top_nodes, name=name_params["set_In"])
247
  outlet_set = my_assembly.Set(nodes=bot_nodes, name=name_params["set_Out"
248
      1)
249
250
  all_left_nodes = my_instance.nodes.getByBoundingBox(
251
       xMin=-1e-3,
252
       yMin = -1e - 3,
253
       zMin = 0.0,
2.54
       xMax=1e-3,
255
      yMax=membrane_thickness + 1e-3,
256
       zMax = 0.0,
257
  )
258
  all_left_nodes_y_coord = []
259
260
  all_right_nodes = my_instance.nodes.getByBoundingBox(
261
       xMin=membrane_width - 1e-3,
262
       yMin = -1e - 3,
263
       zMin = 0.0,
264
       xMax=membrane_width + 1e-3,
265
      yMax=membrane_thickness + 1e-3,
266
       zMax = 0.0,
267
  )
268
  all_right_nodes_y_coord = []
269
270
  for node in all_left_nodes:
       all_left_nodes_y_coord.append(node.coordinates[1])
273
  left_idx = np.argsort(all_left_nodes_y_coord)
274
275
276 left_nodes = []
277 for idx in left_idx:
```

```
for node in all_left_nodes:
278
            if node.coordinates[1] == all_left_nodes_y_coord[idx]:
279
                left_nodes.append(node.label)
280
281
  for node in all_right_nodes:
282
       all_right_nodes_y_coord.append(node.coordinates[1])
283
284
  right_idx = np.argsort(all_right_nodes_y_coord)
285
286
  right_nodes = []
287
  for idx in right_idx:
288
       for node in all_right_nodes:
289
            if node.coordinates[1] == all_right_nodes_y_coord[idx]:
290
                right_nodes.append(node.label)
291
292
293
  node_sets = []
294
  for ii in range(len(left_nodes)):
295
       right_name = "NR" + str(ii)
296
       left_name = "NL" + str(ii)
297
298
       my_assembly.SetFromNodeLabels(
299
            nodeLabels =(
300
                (
301
                     name_params["part"] + "-1",
302
                     (left_nodes[ii],),
303
304
                ),
           ),
305
           name=left_name ,
306
       )
307
       my_assembly.SetFromNodeLabels(
308
            nodeLabels =(
309
                (
                     name_params["part"] + "-1",
311
                     (right_nodes[ii],),
312
                ),
313
           ),
314
315
           name=right_name ,
       )
316
       node_sets.append([right_name, left_name])
317
318
  #
319
320 # Timestep Setup
321 #
    _____
```

```
322
  my_model.MassDiffusionStep(
323
       name=name_params["time_step"],
324
       previous="Initial",
325
       response=STEADY_STATE,
326
       amplitude=RAMP,
327
  )
328
  del my_model.historyOutputRequests["H-Output-1"]
329
  my_model.fieldOutputRequests["F-Output-1"].setValues(variables=("MFL", "
330
      CONC", "COORD", "FV"))
  # _____
  # Setup BCs
333
  # _____
334
335
  inlet_region = regionToolset.Region(nodes=top_nodes)
336
  my_model.ConcentrationBC(
       name=name_params["bc_In"],
338
       createStepName=name_params["time_step"],
339
       region=inlet_region ,
340
       fixed = OFF,
341
       distributionType=UNIFORM,
342
       fieldName="",
343
       magnitude = 1.0,
344
       amplitude=UNSET,
345
  )
346
347
  outlet_region = regionToolset.Region(nodes=bot_nodes)
348
  my_model.ConcentrationBC(
349
       name=name_params["bc_Out"],
350
       createStepName=name_params["time_step"],
351
       region=outlet_region,
       fixed = OFF,
353
       distributionType=UNIFORM,
354
       fieldName="",
355
       magnitude = 0.0,
356
       amplitude=UNSET,
357
358
  )
359
  #
     _____
360
  # Setup Fields
361
  #
    _____
362
363
364 field_region = regionToolset.Region(nodes=interface_nodes_set.nodes)
```

## A Codes

```
my_model.Field(
365
       name=name_params["diff_field"],
366
367
       createStepName=name_params["time_step"],
       region=field region,
368
       distributionType=USER_DEFINED,
369
       fieldVariableNum=1,
370
  )
371
372
  #
     _____
373
  # Setup Constraints
374
  #
    _____
375
376
  for ii , node_set in enumerate(node_sets):
377
       right_name = node_set[0].upper()
378
       left_name = node_set[1].upper()
379
       my_model. Equation (
380
            name = "eq" + str(ii),
381
            terms = ((1.0, right_name, 11), (-1.0, left_name, 11)),
382
       )
383
384
  #
385
  # Setup Job
386
  #
     _____
387
388
  my_job = my_mdb. Job(
389
       name=case,
390
       model=name_params["model"],
391
       description="",
392
       type=ANALYSIS,
393
       atTime=None,
394
       waitMinutes = 0,
395
       waitHours = 0,
396
       queue=None,
397
       memory = 90,
398
       memoryUnits=PERCENTAGE,
399
       getMemoryFromAnalysis=True,
400
       explicitPrecision=SINGLE,
401
       nodalOutputPrecision=SINGLE,
402
       echoPrint=OFF,
403
       modelPrint=OFF,
404
       contactPrint=OFF,
405
       historyPrint=OFF,
406
       userSubroutine="",
407
       scratch = "",
408
```

## A Codes

```
resultsFormat=ODB,
409
       numThreadsPerMpiProcess=1,
410
411
       multiprocessingMode=DEFAULT,
       numCpus=2,
412
       numDomains=2,
413
       numGPUs=1,
414
  )
415
416
417 my_job. writeInput()
```

## A.2.2 ABAQUS Result Extraction

```
1 # Imports
2 from odbAccess import *
3 import numpy as np
4 import os
  import csv
  def log(msg):
8
      # Helper-function to print logs to the python cosole, the abaqus
9
     console and save them in a log file
      print(str(msg))
10
      print >> sys.__stdout__ , str(msg)
11
      with open("../Logs/export_log.txt", "a") as log_file:
          log_file . write ( str (msg) )
13
          log_file . write("\n")
14
15
16
 case = None
17
  with open("../DATA_FILES/actual_case.txt") as file:
18
      reader = csv.reader(file, delimiter=",")
19
      for row in reader:
20
          if row:
               case = row[0].replace("\\", "/").replace("DATA_FILES/", "")
22
23
24 case_path = ".../DATA_FILES/" + case
25
26 name_params = None
  with open(case_path + "/Output_Data.csv") as file:
27
      reader = csv.DictReader(file, delimiter=",")
28
      for row in reader:
29
```
```
name_params = row
30
31
32 # Export Data
my_session = session
34 my_odb = my_session.openOdb(name=case + ".odb")
35 frame = my_odb.steps[name_params["time_step"]].frames[1]
<sup>36</sup> instance = my_odb.rootAssembly.instances[name_params["part"].upper() + "
     -1"]
37
  set_region = my_odb.rootAssembly.nodeSets[name_params["set_Out"].upper()
38
     1
  mfl_data = frame.fieldOutputs["MFL"].getSubset(region=set_region,
39
     position = ELEMENT_NODAL).values
40
  data_list = []
41
  for index in range(len(mfl_data)):
42
      data_list.append(
43
          [
44
               instance.getNodeFromLabel(mfl_data[index].nodeLabel).
45
     coordinates [0],
               instance.getNodeFromLabel(mfl_data[index].nodeLabel).
46
     coordinates [1],
               mfl_data[index].data[1],
47
          ]
48
      )
49
50 np.savetxt(
      "../RESULTS/Result_" + case + ".csv",
51
      data_list,
52
      delimiter=", ",
53
      fmt="% s",
54
      header="x, y, mfl_2",
55
56)
```

## A.2.3 ABAQUS Standard Simulation

```
# Imports
2 from abagus import *
  from abaqusConstants import *
3
4 from caeModules import *
6 import csv
 import math
8 import numpy as np
9 import re
10 import sys
13 def log(msg):
      # Helper-function to print logs to the python cosole, the abaqus
14
     console and save them in a log file
      print(str(msg))
15
      print >> sys.__stdout__ , str(msg)
16
      with open("../Logs/log.txt", "a") as log_file:
17
          log_file . write ( str (msg) )
18
          log_file . write ( "\n" )
19
20
22 # Variables
_{23} case = None
  with open("../DATA_FILES/actual_case.txt") as file:
24
      reader = csv.reader(file, delimiter=",")
25
      for row in reader:
26
           if row:
27
               case = row[0].replace("\\", "/").replace("DATA_FILES/", "")
28
29
30
31 case_path = ".../DATA_FILES/" + case
32
33 geometric_params = []
  with open(case_path + "/Geometric_Data.csv") as file:
34
      reader = csv.reader(file, delimiter=",")
35
      for row in reader:
36
           for row_elem in row:
37
               str_floats = re.findall(r"[-+]?\d*\.\d+l\d+", row_elem)
38
               tmp = []
39
               for str_elem in str_floats:
40
                   tmp.append(float(str_elem))
41
```

```
geometric_params.append(tmp)
42
43
44 name_params = None
  with open(case path + "/Output Data.csv") as file:
45
      reader = csv.DictReader(file, delimiter=",")
46
      for row in reader:
47
          name_params = row
48
49
50 filler_radius = geometric_params[0][0]
st interface_thickness = geometric_params[1][0]
52 membrane_thickness = geometric_params [2][0]
  membrane_width = geometric_params[3][0]
53
54
solubility = geometric_params[4][0]
56 diffusivity_matrix = geometric_params[5][0]
57 diffusivity_interface = geometric_params[6][0]
s8 interface_prefactor = geometric_params[7][0]
59
60 matrix_mesh_size = geometric_params[8][0]
  interface_mesh_size = geometric_params[9][0]
61
62
  filler_center = [geometric_params[10][0], geometric_params[11][0]]
63
64
65 # _____
 # Abaqus Setup
66
 # _____
67
68
69 Mdb()
70 \text{ my}_{session} = session
71 \text{ my} \text{mdb} = \text{mdb}
72
<sup>73</sup> my_mdb.models.changeKey(fromName="Model-1", toName=name_params["model"])
74 my_model = my_mdb.models[name_params["model"]]
75
 # _____
76
 # Part Creation
77
 # _____
78
79
sol sketch = my_model.ConstrainedSketch(name="__profile__", sheetSize=200.0)
sketch.rectangle(point1=(0, 0), point2=(membrane_width),
     membrane_thickness))
82 p1 = (filler_center[0] + filler_radius, filler_center[1])
83 sketch.CircleByCenterPerimeter(center=filler_center, point1=p1)
```

```
my_part = my_model. Part (name=name_params ["part"], dimensionality=
84
     TWO_D_PLANAR, type=DEFORMABLE_BODY)
85
  my_part.BaseShell(sketch=sketch)
86
  sketch = my_model.ConstrainedSketch(name="__profile__", sheetSize=200.0)
87
  p2 = (filler_center[0] + filler_radius + interface_thickness,
88
      filler_center [1])
  sketch.CircleByCenterPerimeter(center=filler_center, point1=p2)
89
  my_part.PartitionFaceBySketch(faces=my_part.faces, sketch=sketch)
90
91
  # _____
92
  # Material Setup
93
   _____
  #
94
95
% my_material = my_model.Material(name=name_params["material_matrix"])
  my_material.Diffusivity(law=GENERAL, table=((diffusivity_matrix,),))
97
  my_material. Solubility (table = (( solubility ,) ,) )
98
99
  # _____
100
  # Section Setup
101
  # _____
102
103
  my_model.HomogeneousSolidSection(
104
      name=name_params["section_matrix"], material=name_params["
105
      material_matrix"], thickness=None
  )
106
107
  region_matrix = regionToolset.Region(faces=my_part.faces)
108
109
  my_part.SectionAssignment(
110
      region=region_matrix,
111
      sectionName=name_params["section_matrix"],
      offset = 0.0,
113
      offsetType=MIDDLE_SURFACE,
114
      offsetField="",
115
      thicknessAssignment=FROM_SECTION,
116
117)
118
  # _____
119
120 # Mesh Setup
  #
    _____
121
elemType1 = mesh.ElemType(elemCode=DC2D4, elemLibrary=STANDARD)
elemType2 = mesh.ElemType(elemCode=DC2D3, elemLibrary=STANDARD)
```

```
125
  my_part.setMeshControls(regions=my_part.faces, elemShape=QUAD, algorithm
126
      =ADVANCING_FRONT)
  my part.setElementType(
127
      regions=regionToolset.Region(faces=my_part.faces),
128
       elemTypes=(elemType1, elemType2),
129
  )
130
  my_part.seedPart(size=matrix_mesh_size, deviationFactor=0.1,
      minSizeFactor = 0.1)
  interface_edges = my_part.edges.getByBoundingBox(
133
      xMin=(filler_center[0] - filler_radius - interface_thickness - 1e-3)
134
      xMax=(filler_center[0] + filler_radius + interface_thickness + 1e-3)
135
      yMin=(filler_center[1] - filler_radius - interface_thickness - 1e-3)
136
      yMax=(filler_center[1] + filler_radius + interface_thickness + 1e-3)
138)
139 my_part.Set(edges=interface_edges, name="SeedByEdge")
140 my_part.seedEdgeBySize(edges=interface_edges, size=interface_mesh_size)
  my_part.generateMesh()
141
142
  # _____
143
144 # Assembly Setup
  # _____
145
146
  my_assembly = my_model.rootAssembly
147
  my_assembly.DatumCsysByDefault(CARTESIAN)
148
149
  my_{instance} = my_{assembly}. Instance (name=name_params["part"] + "-1", part
150
      =my_part, dependent=ON)
151
152
  # _____
153 # Set Setup
  #
154
155
  all_nodes = my_instance.nodes
156
  top_nodes = my_instance.nodes.getByBoundingBox(
157
      xMin = -1e - 3,
158
      yMin=membrane_thickness - 1e-3,
159
      zMin = 0.0,
160
      xMax=membrane_width + 1e-3,
161
```

```
yMax=membrane_thickness + 1e-3,
162
      zMax = 0.0,
  )
164
  bot nodes = my instance.nodes.getByBoundingBox(
165
      xMin=-1e-3, yMin=-1e-3, zMin=0.0, xMax=membrane_width + 1e-3, yMax=1
166
      e - 3, zMax = 0.0
  )
167
168
  interface_nodes = my_instance.nodes.getByBoundingCylinder(
169
       center1 = (filler_center[0], filler_center[1], 0 - 1e-3),
170
       center2 = (filler_center[0], filler_center[1], 0 + 1e-3),
       radius=filler_radius + interface_thickness + 1e-3,
  )
174
  interface_elements = my_instance.elements.getByBoundingCylinder(
       center1 = (filler_center[0], filler_center[1], 0 - 1e-3),
176
       center2 = (filler_center[0], filler_center[1], 0 + 1e-3),
       radius=filler_radius + interface_thickness + 1e-3,
178
  )
179
180
  all_nodes_set = my_assembly.Set(nodes=all_nodes, name=name_params["
181
      set_All"])
  interface_nodes_set = my_assembly.Set(nodes=interface_nodes, name=
182
      name_params["set_Interface_Nodes"])
  interface_element_set = my_assembly.Set(nodes=interface_elements, name=
183
      name_params["set_Interface_Elements"])
  inlet_set = my_assembly.Set(nodes=top_nodes, name=name_params["set_In"])
184
  outlet_set = my_assembly.Set(nodes=bot_nodes, name=name_params["set_Out"
185
      ])
186
187
  all_left_nodes = my_instance.nodes.getByBoundingBox(
188
      xMin = -1e - 3,
189
      yMin = -1e - 3,
190
      zMin = 0.0,
191
      xMax=1e-3,
192
      yMax=membrane_thickness + 1e-3,
193
      zMax = 0.0,
194
  )
195
  all_left_nodes_y_coord = []
196
197
  all_right_nodes = my_instance.nodes.getByBoundingBox(
198
      xMin=membrane_width - 1e-3,
199
      yMin=-1e-3,
200
```

```
zMin = 0.0,
201
       xMax=membrane_width + 1e-3,
202
203
       yMax=membrane_thickness + 1e-3,
       zMax = 0.0,
204
  )
205
  all_right_nodes_y_coord = []
206
207
  for node in all_left_nodes:
208
       all_left_nodes_y_coord.append(node.coordinates[1])
209
210
  left_idx = np.argsort(all_left_nodes_y_coord)
212
  left_nodes = []
213
  for idx in left_idx:
214
       for node in all_left_nodes:
215
           if node.coordinates[1] == all_left_nodes_y_coord[idx]:
216
                left_nodes.append(node.label)
218
  for node in all_right_nodes:
219
       all_right_nodes_y_coord.append(node.coordinates[1])
220
  right_idx = np.argsort(all_right_nodes_y_coord)
224 right_nodes = []
  for idx in right_idx:
225
       for node in all_right_nodes:
226
           if node.coordinates[1] == all_right_nodes_y_coord[idx]:
227
                right_nodes.append(node.label)
228
229
230
  node_sets = []
  for ii in range(len(left_nodes)):
       right name = "NR" + str(ii)
       left_name = "NL" + str(ii)
234
235
       my_assembly.SetFromNodeLabels(
236
           nodeLabels =(
                (
238
                    name_params["part"] + "-1",
239
                    (left_nodes[ii],),
240
                ),
241
           ),
242
           name=left_name ,
243
       )
244
```

```
my_assembly.SetFromNodeLabels(
245
            nodeLabels =(
246
24
                (
                     name_params ["part"] + "-1",
248
                     (right_nodes[ii],),
249
                ),
250
            ),
251
            name=right_name ,
2.52
       )
253
       node_sets.append([right_name, left_name])
254
255
  #
256
    Timestep Setup
  #
257
  #
258
2.59
  my_model.MassDiffusionStep(
260
       name=name_params["time_step"],
261
       previous="Initial",
262
       response=STEADY_STATE,
263
       amplitude=RAMP,
264
  )
265
  del my_model. historyOutputRequests ["H-Output-1"]
266
  my_model.fieldOutputRequests["F-Output-1"].setValues(variables=("MFL", "
267
      CONC", "COORD", "FV"))
268
  #
    _____
269
  # Setup BCs
270
  #
    _____
272
  inlet_region = regionToolset.Region(nodes=top_nodes)
273
  my_model.ConcentrationBC(
274
       name=name_params["bc_In"],
275
       createStepName=name_params["time_step"],
276
       region=inlet_region ,
277
       fixed = OFF,
278
       distributionType=UNIFORM,
279
       fieldName="",
280
       magnitude = 1.0,
281
       amplitude=UNSET,
282
  )
283
284
  outlet_region = regionToolset.Region(nodes=bot_nodes)
285
  my_model.ConcentrationBC(
286
       name=name_params["bc_Out"],
287
```

```
createStepName=name_params["time_step"],
288
       region=outlet_region,
289
       fixed = OFF,
290
       distributionType=UNIFORM,
291
       fieldName="",
292
       magnitude = 0.0,
293
       amplitude=UNSET,
294
  )
295
296
  #
297
  #
    Setup Constraints
298
  #
       ____
299
300
  for ii, node_set in enumerate(node_sets):
301
       right_name = node_set[0].upper()
302
       left_name = node_set[1].upper()
303
       my_model.Equation(
304
            name="eq" + str(ii),
305
            terms = ((1.0, right_name, 11), (-1.0, left_name, 11)),
306
307
       )
308
309
  #
  # Setup Job
311
  #
      _____
313
  job_name = "Standart__" + case
314
  my_job = my_mdb. Job(
315
       name=job_name ,
316
       model=name_params["model"],
317
       description="",
318
       type=ANALYSIS,
       atTime=None,
320
       waitMinutes = 0,
321
       waitHours=0,
       queue=None,
323
       memory = 90,
324
       memoryUnits=PERCENTAGE,
325
       getMemoryFromAnalysis=True,
326
       explicitPrecision=SINGLE,
       nodalOutputPrecision=SINGLE,
       echoPrint=OFF,
329
       modelPrint=OFF.
330
       contactPrint=OFF,
331
```

```
historyPrint=OFF,
332
       userSubroutine="",
333
334
       scratch="",
       resultsFormat=ODB,
335
       numThreadsPerMpiProcess=1,
336
       multiprocessingMode=DEFAULT,
       numCpus=2,
338
       numDomains=2,
339
      numGPUs=1,
340
  )
341
342
  my_job.submit()
343
  my_job.waitForCompletion()
344
345
  # _____
346
  # Output
347
  # _____
348
349
my_session = session
351 my_odb = my_session.openOdb(name=job_name + ".odb")
352 frame = my_odb.steps[name_params["time_step"]].frames[1]
instance = my_odb.rootAssembly.instances[name_params["part"].upper() + "
      -1"]
354
  set_region = my_odb.rootAssembly.nodeSets[name_params["set_Out"].upper()
355
      ]
  mfl_data = frame.fieldOutputs ["MFL"].getSubset (region=set_region,
356
      position = ELEMENT_NODAL). values
357
  data_list = []
358
  for index in range(len(mfl_data)):
359
       data_list.append(
360
           ſ
361
                instance.getNodeFromLabel(mfl_data[index].nodeLabel).
362
      coordinates [0],
                instance.getNodeFromLabel(mfl_data[index].nodeLabel).
363
      coordinates [1],
                mfl_data[index].data[1],
364
           ]
365
       )
366
  np.savetxt(
367
       "../RESULTS/Result_" + job_name + ".csv",
368
       data_list,
369
       delimiter=", ",
370
```

```
371 fmt="% s",
372 header="x, y, mfl_2",
373 )
```

## A.2.4 Postprocessing

```
# Imports
 import os
3 import numpy as np
4 import matplotlib.pyplot as plt
  import scipy.integrate as sci
 plt.rc("font", size=18) # controls default text sizes
7
 plt.rc("legend", fontsize=10) # legend fontsize
8
 plt.rc("figure", titlesize=18) # fontsize of the figure title
 plt.rc("axes", labelsize=18) # fontsize of the x and y labels
10
12 # PCCL-Colors
<sup>13</sup> PCCL_main = tuple(np. array([0.0, 125.0, 166.0]) / 255.0)
<sup>14</sup> PCCL_lightblue = tuple (np. array ([1.0, 174.0, 240.0]) / 255.0)
<sup>15</sup> PCCL_darkgreen = tuple(np. array([56.0, 181.0, 77.0]) / 255.0)
16 PCCL_lightgreen = tuple (np. array ([153.0, 202.0, 59.0]) / 255.0)
17 PCCL_yellow = tuple (np. array ([232.0, 196.0, 56.0]) / 255.0)
<sup>18</sup> PCCL_orange = tuple (np. array ([238.0, 138.0, 44.0]) / 255.0)
<sup>19</sup> PCCL_red = tuple (np. array ([230.0, 47.0, 41.0]) / 255.0)
_{20} PCCL_purple = tuple (np. array ([131.0, 42.0, 132.0]) / 255.0)
21 PCCL_violet = tuple (np. array ([92.0, 73.0, 152.0]) / 255.0)
22 PCCL_darkblue = tuple (np. array ([0.0, 102.0, 176.0]) / 255.0)
23
 pccl_colors = [
24
      PCCL_main,
25
      PCCL_red,
26
      PCCL_darkgreen,
27
      PCCL_orange,
28
      PCCL_purple,
29
      PCCL_darkblue,
30
31
  1
_{33} linestyles = {
      "loosely dotted": (0, (1, 10)),
34
      "dotted": (0, (1, 1)),
35
      "densely dotted": (0, (1, 1)),
36
```

```
"long dash with offset": (5, (10, 3)),
37
      "loosely dashed": (0, (5, 10)),
38
      "dashed": (0, (5, 5)),
39
      "densely dashed": (0, (5, 1)),
40
      "loosely dashdotted": (0, (3, 10, 1, 10)),
41
      "dashdotted": (0, (3, 5, 1, 5)),
42
      "densely dashdotted": (0, (3, 1, 1, 1)),
43
      "dashdotdotted": (0, (3, 5, 1, 5, 1, 5)),
44
      "loosely dashdotdotted": (0, (3, 10, 1, 10, 1, 10)),
45
      "densely dashdotdotted": (0, (3, 1, 1, 1, 1, 1)),
46
47
 }
48
49 result_directory = "RESULTS/"
50 file_names = []
  for file in os.listdir(result_directory):
51
      if file.endswith(".csv"):
52
          if "Outflow" not in file:
53
              file_names.append(file)
54
56 data_arr = []
 for file_name in file_names:
57
      data_arr.append(np.genfromtxt(result_directory + file_name,
58
     delimiter=","))
59
60 plot_list = []
  outflow_data = []
61
  for idx, data in enumerate(data_arr):
62
      pos_arr = np.array([item[0] for item in data])
63
      mfl_2_arr = np. array ([item [2] for item in data])
64
64
      sorted_pos_arr, sort_idx = np.unique(pos_arr, return_index=True)
66
      sorted_mfl_2_arr = np.zeros(sorted_pos_arr.size)
67
      for unique_idx , unique_pos in enumerate(sorted_pos_arr):
68
          value_idx = np.where(pos_arr == unique_pos)
69
          sorted_mfl_2_arr[unique_idx] = np.mean(mfl_2_arr[value_idx])
70
71
      integrated_value = sci.simpson(sorted_mfl_2_arr, sorted_pos_arr)
72
      membrane_length = sorted_pos_arr[-1] - sorted_pos_arr[0]
73
      outflow_value = integrated_value / membrane_length * -1.0
74
75
      plot_label = (
76
          file_names[idx]
77
          .replace(".csv", "")
78
          .replace("Result_Job__", "")
79
```

```
.replace("FS_5-0_2-0_IT_0-5_DM_1-0_", "")
80
           .replace("PF_", "psi: ")
81
           .replace("DI_", "D: ")
82
           .replace("-", ".")
83
           .replace("__", ", ")
84
       )
85
       if "Standart" in plot_label:
86
           plot_label = "No Interface"
87
88
       outflow_data.append([plot_label, outflow_value])
89
       plot_list.append([plot_label, sorted_pos_arr, sorted_mfl_2_arr])
90
91
  plt.figure()
92
  plt.xlabel("Membrane Width")
93
  plt.ylabel("Outflow over Membrane Width")
94
95
  for plot_data in plot_list:
96
       plot_label = plot_data[0]
97
       plot_color = pccl_colors[1]
98
       plot_linestyle = "-"
99
100
       if "psi: 75" in plot_label:
101
           plot_linestyle = linestyles["dashed"]
       elif "psi: 50" in plot_label:
103
           plot_linestyle = linestyles["dashdotted"]
104
       elif "psi: 20" in plot_label:
105
           plot_linestyle = linestyles["dotted"]
106
       elif "psi: 5" in plot_label:
107
           plot_linestyle = linestyles["dashdotdotted"]
108
109
       if "D: 10.0," in plot_label:
110
           plot_color = pccl_colors[0]
       elif "D: 1.0," in plot_label:
           plot_color = pccl_colors[2]
       elif "D: 0.1," in plot_label:
114
           plot_color = pccl_colors[3]
115
       elif "D: 0.01," in plot_label:
116
           plot_color = pccl_colors[4]
       elif "D: 0.001," in plot_label:
118
           plot_color = pccl_colors[5]
119
120
       plt.plot(
           plot_data[1],
           plot_data[2],
123
```

```
label=plot_label ,
124
           color=plot_color ,
125
           linestyle=plot_linestyle,
126
       )
127
128
  plt.legend()
129
  # plt.tight_layout()
130
  plt.figure()
132
  plt.xlabel("Membrane Width")
  plt.ylabel("Outflow over Membrane Width")
134
  for plot_data in plot_list:
135
       if "No Interface" in plot_data[0] or "psi: 20" in plot_data[0]:
136
           plot_label = plot_data[0]
           plot_color = pccl_colors[1]
138
           plot_linestyle = "-"
139
140
           if "psi: 75" in plot_label:
141
                plot_linestyle = linestyles["dashed"]
142
           elif "psi: 50" in plot_label:
143
                plot_linestyle = linestyles["dashdotted"]
144
           elif "psi: 20" in plot_label:
145
                plot_linestyle = linestyles["dotted"]
146
           elif "psi: 5" in plot_label:
147
                plot_linestyle = linestyles["dashdotdotted"]
148
149
           if "D: 10.0," in plot_label:
150
                plot_color = pccl_colors[0]
151
           elif "D: 1.0," in plot_label:
                plot_color = pccl_colors[2]
153
           elif "D: 0.1," in plot_label:
154
                plot_color = pccl_colors[3]
155
           elif "D: 0.01," in plot_label:
156
                plot_color = pccl_colors[4]
157
           elif "D: 0.001," in plot_label:
158
                plot_color = pccl_colors[5]
159
160
           plt.plot(
161
                plot_data[1],
162
                plot_data[2],
163
                label=plot_label ,
164
                color=plot_color ,
165
                linestyle=plot_linestyle,
166
           )
167
```

```
168
  plt.legend()
169
170 # plt.tight_layout()
  plt.savefig("Outflows_Rect_Small_New.png", bbox_inches="tight")
171
173
  plt.figure()
174
175
  x_{positions} = [1, 2, 3, 4]
176
|x_axis_labels = ["5", "20", "50", "75"]
178 plt. xlabel ("psi")
  plt.xticks(x_positions, x_axis_labels)
179
180
  plt.ylabel("Sum of Outflow over Membrane Width")
181
182
  for elem in outflow_data:
183
       plot_label = elem[0]
184
       mark_style = "o"
185
       plot_color = pccl_colors[1]
186
       x_{pos} = 0
187
188
       if "psi: 75" in plot_label:
189
           x_pos = 4
190
       elif "psi: 50" in plot_label:
191
           x_pos = 3
192
       elif "psi: 20" in plot_label:
193
           x_pos = 2
194
       elif "psi: 5" in plot_label:
195
           x_{pos} = 1
196
197
       if "D: 10.0," in plot_label:
198
           plot_color = pccl_colors[0]
199
       elif "D: 1.0," in plot_label:
200
           plot_color = pccl_colors[2]
201
       elif "D: 0.1," in plot_label:
202
           plot_color = pccl_colors[3]
203
       elif "D: 0.01," in plot_label:
204
           plot_color = pccl_colors[4]
205
       elif "D: 0.001," in plot_label:
206
           plot_color = pccl_colors[5]
207
208
       if not "No Interface" in plot_label:
209
           plt.plot(x_pos, elem[1], color=plot_color, markersize=4, label=
      plot_label , marker=mark_style , linewidth=0)
```

```
211 else:
212 plt.hlines(xmin=1, xmax=4, y=elem[1], color=plot_color, label="
No Interface")
213
214
215 plt.legend()
# plt.tight_layout()
217 plt.show()
```